

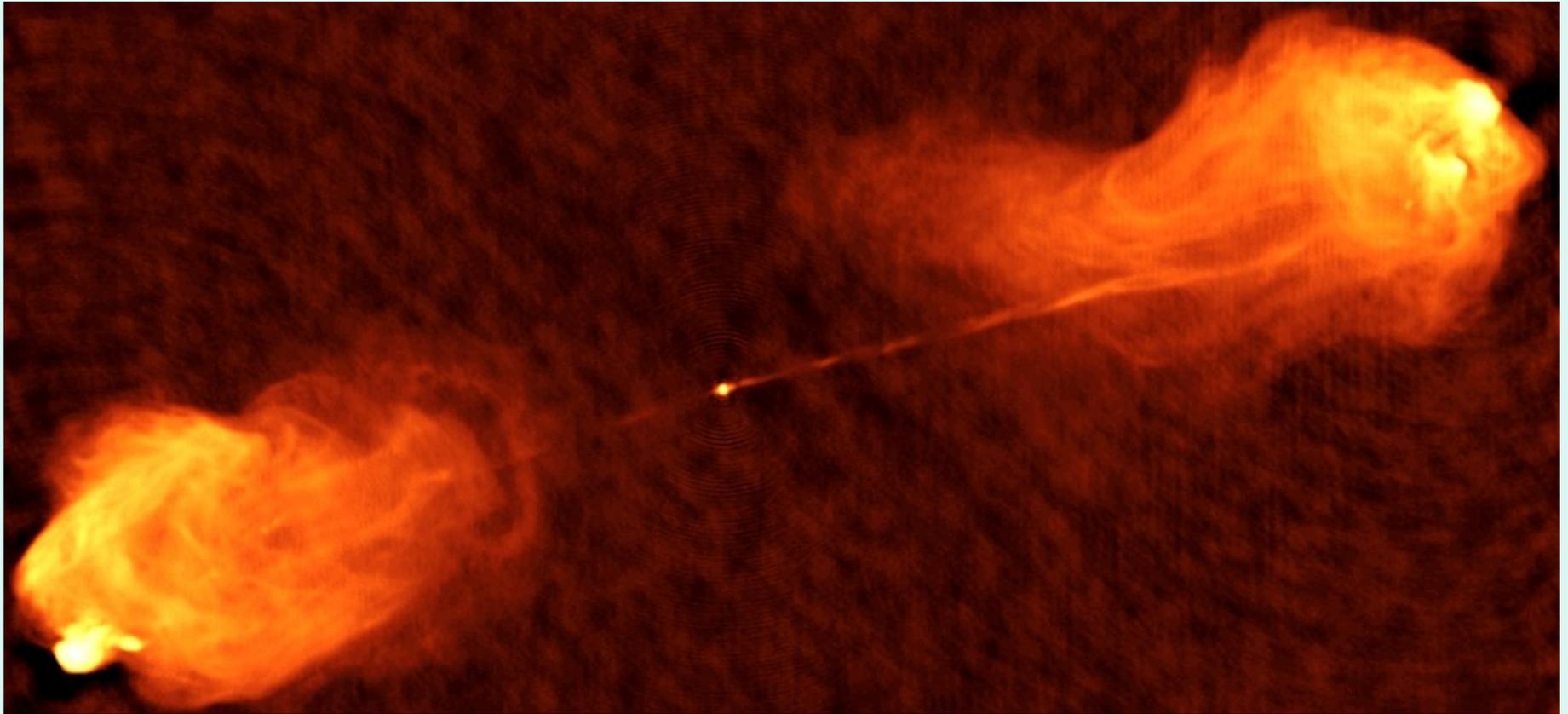
Квазар = QSO = quasi stellar object (1963, Maarten Schmidt)

Активные галактические ядра = квазары $\times 0.01 - 0.00001$

Тепловая светимость $10^{40} - 40^{47}$ эрг/с

Суть явления: на черную дыру массой $10^6 - 10^{10} = M_{\odot}$ падает вещество.

Аккреционные диски – источник света и джетов



FR II

Cyg A FR II

CYGNUS A

VLA 6 cm

VLBI 18 cm

VLBI 1.3 cm

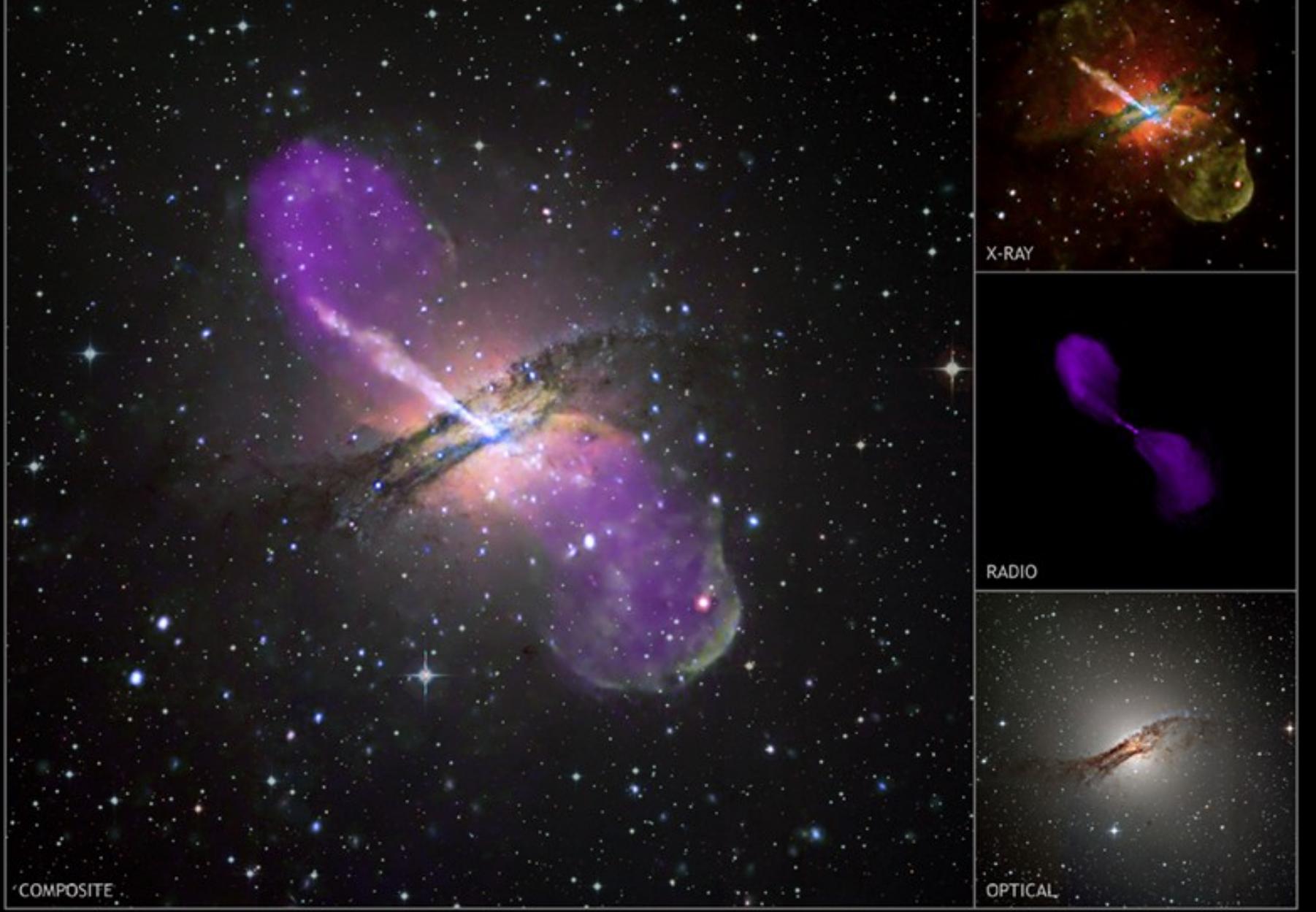
VLBI 7 mm

Lyrs
37500

Lyrs
270

Lyrs
10

Lyrs
4

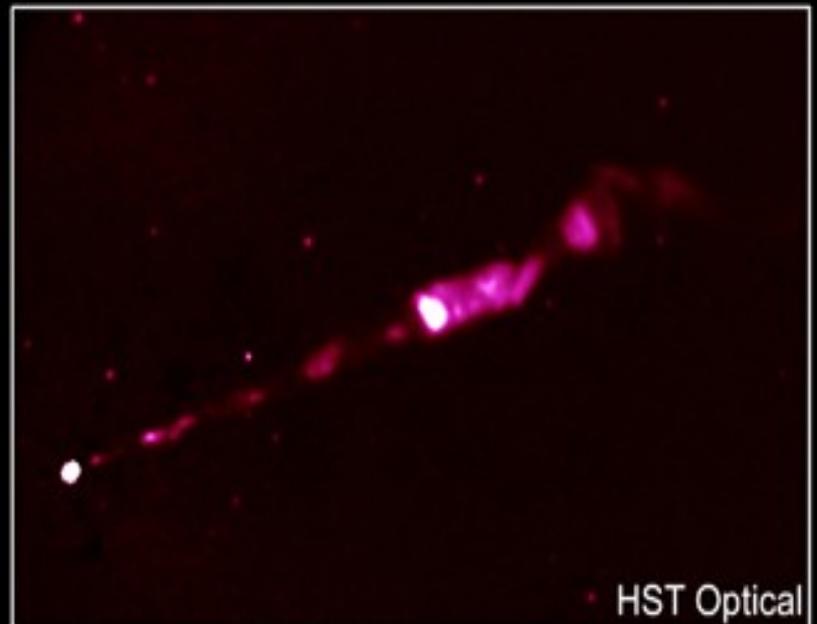
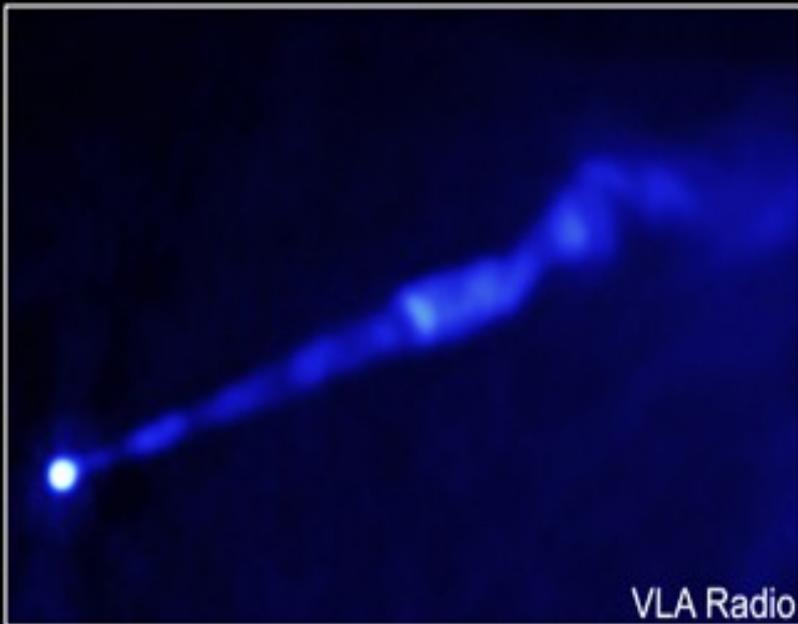
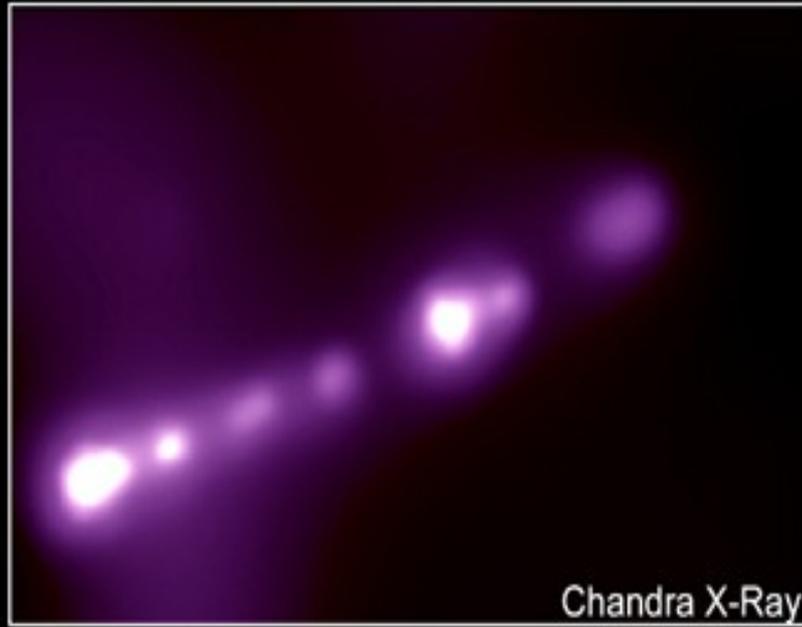


Cen A FR I

TeV emission

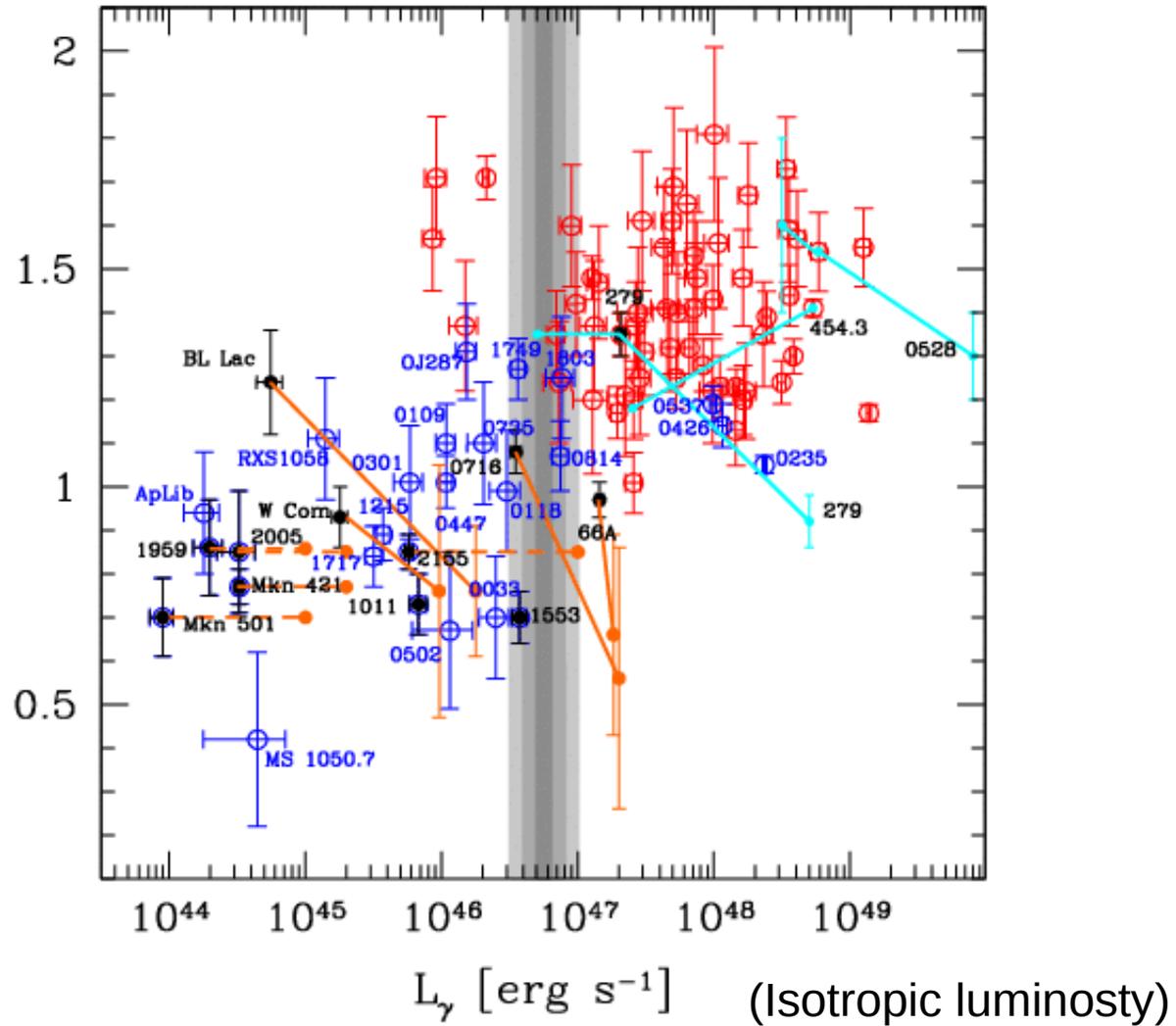
At 20 degrees

The external
environment or
a slow sheath?



Blazars/BL-Lacs Fermi data

(From Ghisellini 2009)



Average power law index in Fermi range

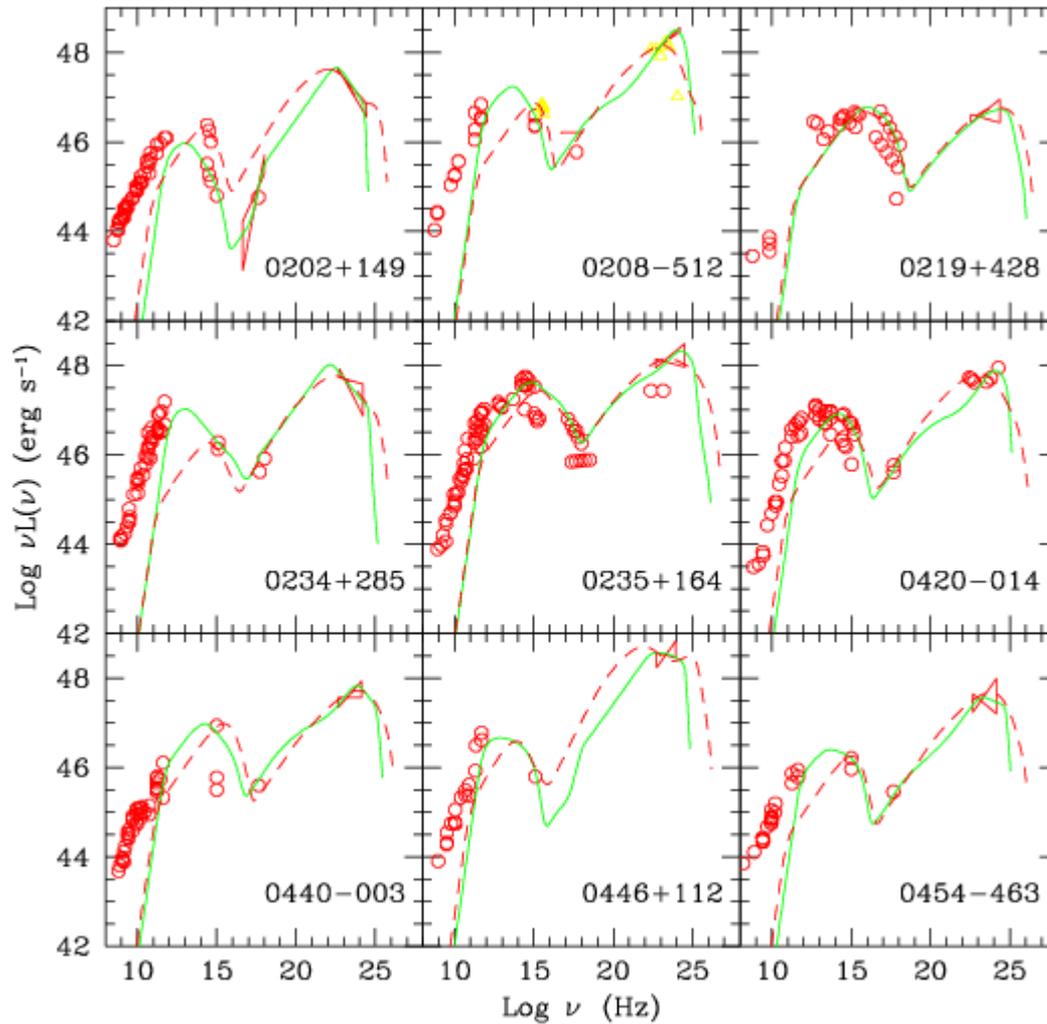
Spectra

Fit:

Synchrotron
-self
Compton

+

External
Compton

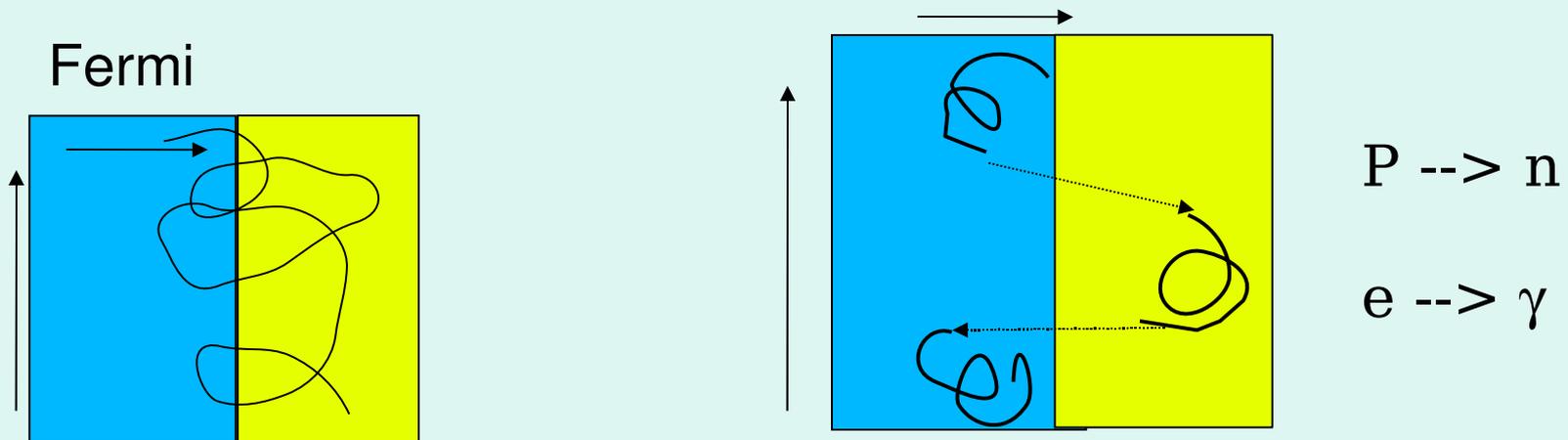


Converter mechanism: an alternative way of energy transfer from a relativistic fluid to particles

Fermi acceleration: a charged particle crosses a shock front back and forth gaining factor Γ^2 in energy at each round trip.

It is not efficient in the case of relativistic shocks.

Derishev et al. (2003) and Stern (2003) independently suggested an alternative: **converter mechanism**.



Charge conversion:

1. **proton** + **photon** \rightarrow **n** + **π^+** (photomeson production)

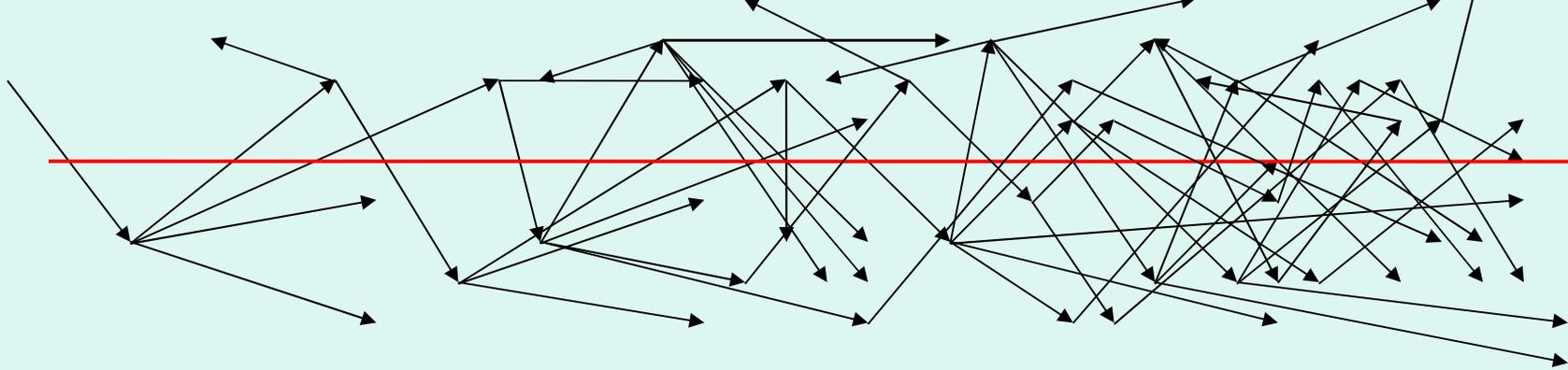
2. **electron** + **photon** \rightarrow **electron** + **photon**
(inverse Compton scattering)

3. **photon** + **photon** \rightarrow **e^+** **e^-**

A background of **low energy photons** with column density $> 10^{25}/\text{cm}^2$ is required

Fermi mechanism: small energy losses and a slow energy gain at a large Γ

Converter mechanism: large energy losses at the charge exchange and a fast energy gain at a large Γ



Photon breeding is the version of the converter mechanism with $\gamma \rightarrow e^+e^- \rightarrow \gamma$ conversion in the supercritical regime (nuclear pile).

Breeding rather than **acceleration** because the spectrum of particles evolves slowly while their number rapidly increases.

Stern & Poutanen (2006, 2008)

Dissipation of the jet bulk energy into gamma-rays via
the photon breeding

Photon breeding mechanism in jets and its observational signatures

Stern & Poutanen 2006 MNRAS (description of the mechanism, analytical estimates problem setup, 1D simulations)

Stern & Poutanen 2008 MNRAS (more extensive 2D simulations, astrophysical implications)

Stern & Tikhomirova 2009 MNRAS accepted (possible observational signature: synchrotron cutoff in GeV range)

The jet



C3
(Comptonization)

C2 $\sim \Gamma^2$

C1

C4 (escape from the jet + pair production)

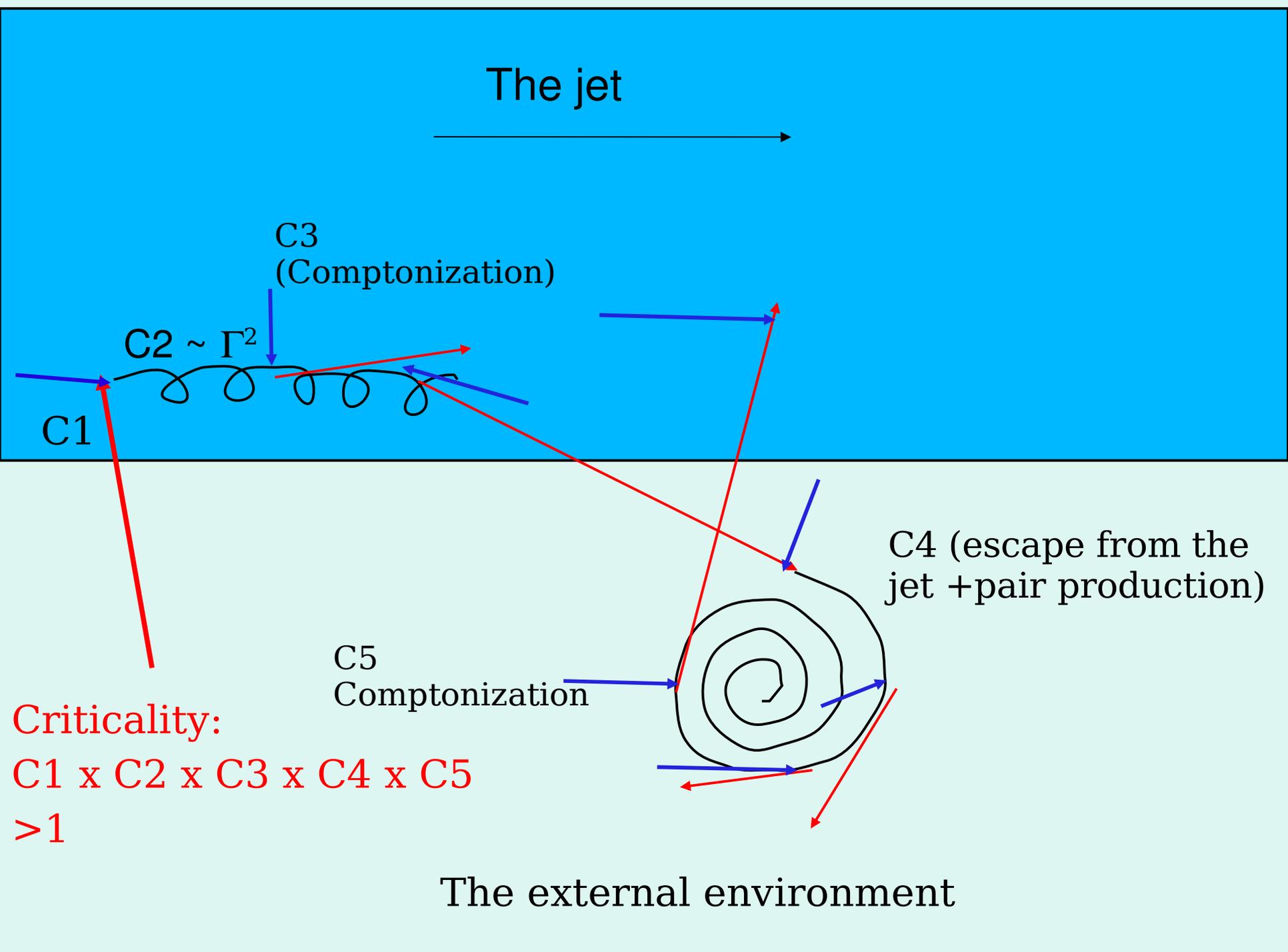
C5
Comptonization

Criticality:

$C1 \times C2 \times C3 \times C4 \times C5$

> 1

The external environment



What we need to launch the mechanism:

1. Transversal magnetic field in the jet H_j --> energy gain.
2. Transversal magnetic field H_e outside --> reflection.
3. Soft background radiation. --> charge conversion

(i) Radiation of the disk, F_d , (mainly along the jet) - insufficient

(ii) A scattered, nearly isotropic component F_t .

6. A seed high energy photon.

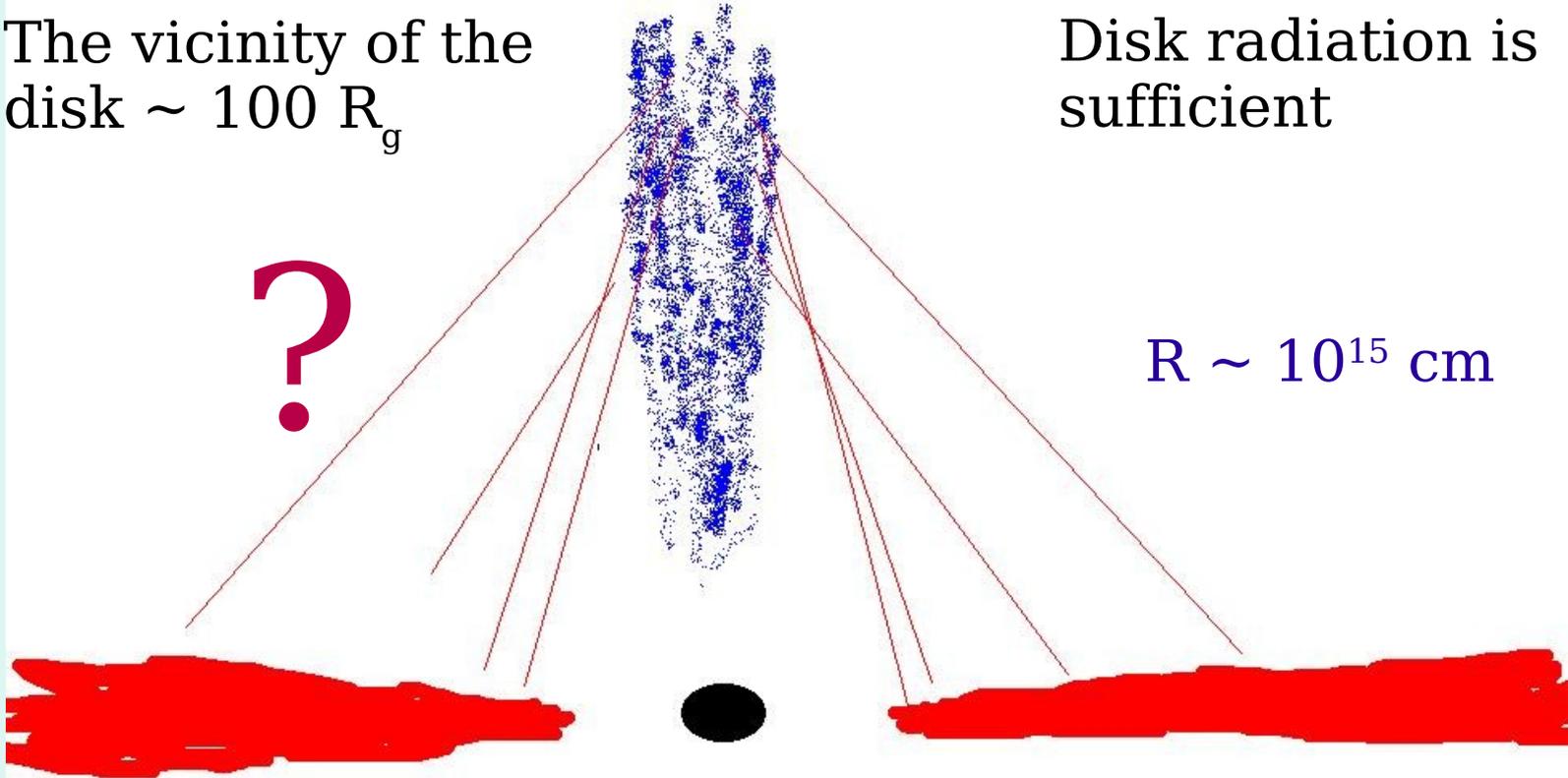
Photon breeding sites

The vicinity of the
disk $\sim 100 R_g$

Disk radiation is
sufficient

?

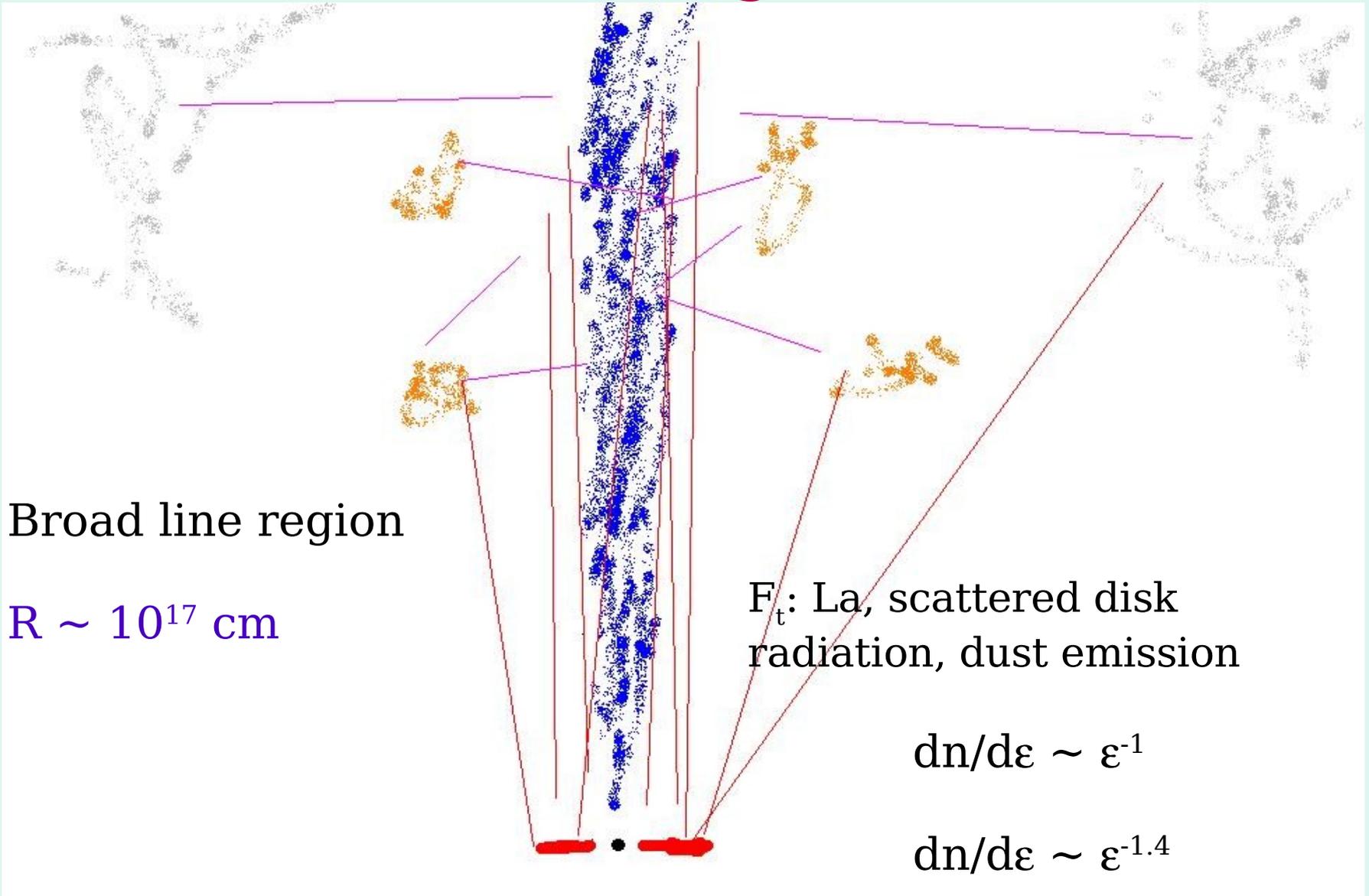
$R \sim 10^{15}$ cm



Model C

No F_t

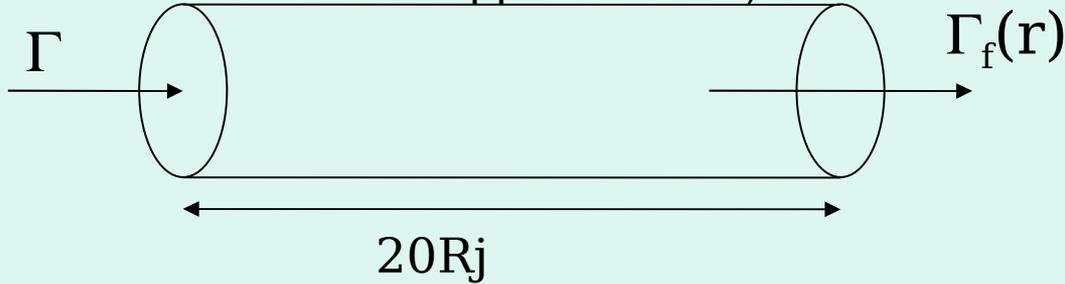
Photon breeding sites



The numerical model

Jet: a cylinder or a cone split into 64 X 20 cells which can independently decelerate, (two

dimensional ballistic approximation). $R = 2 \cdot 10^{17}$ cm, $R_j = 10^{16}$ cm



Soft photon field: disk radiation (multicolor spectrum) along the jet + isotropic radiation, spectral number index $\Gamma = -1.4$, $x_{\max} = T_{\max} = 5$ eV

$$F_{\text{iso}} = F_d / 20$$

Magnetic field:

$$(B_{\text{ext}} = 10^{-3} \text{ G}).$$

Boundary conditions: $\Gamma = \text{const}$ at the inlet The fixed cylindrical boundary. The external

Numerical experiments with a relativistic radiating jet

The tool: Large Particle Monte-Carlo Code (nonlinear) – LPMC

Stern et al 1995

Synchrotron radiation

Compton scattering

Photon-photon – e⁺e⁻ pair production

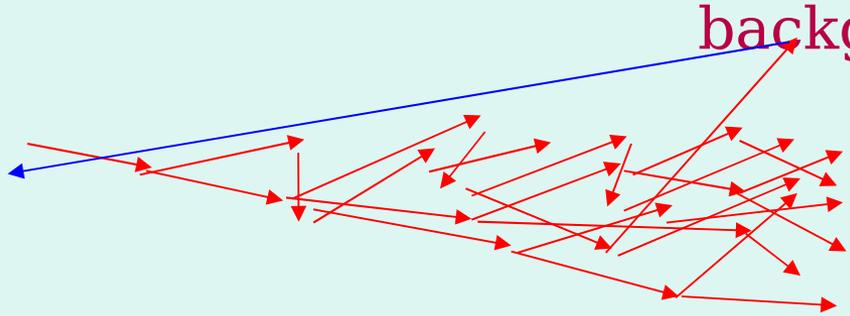
Pair annihilation

Synchrotron selfabsorption

Bremstahlung

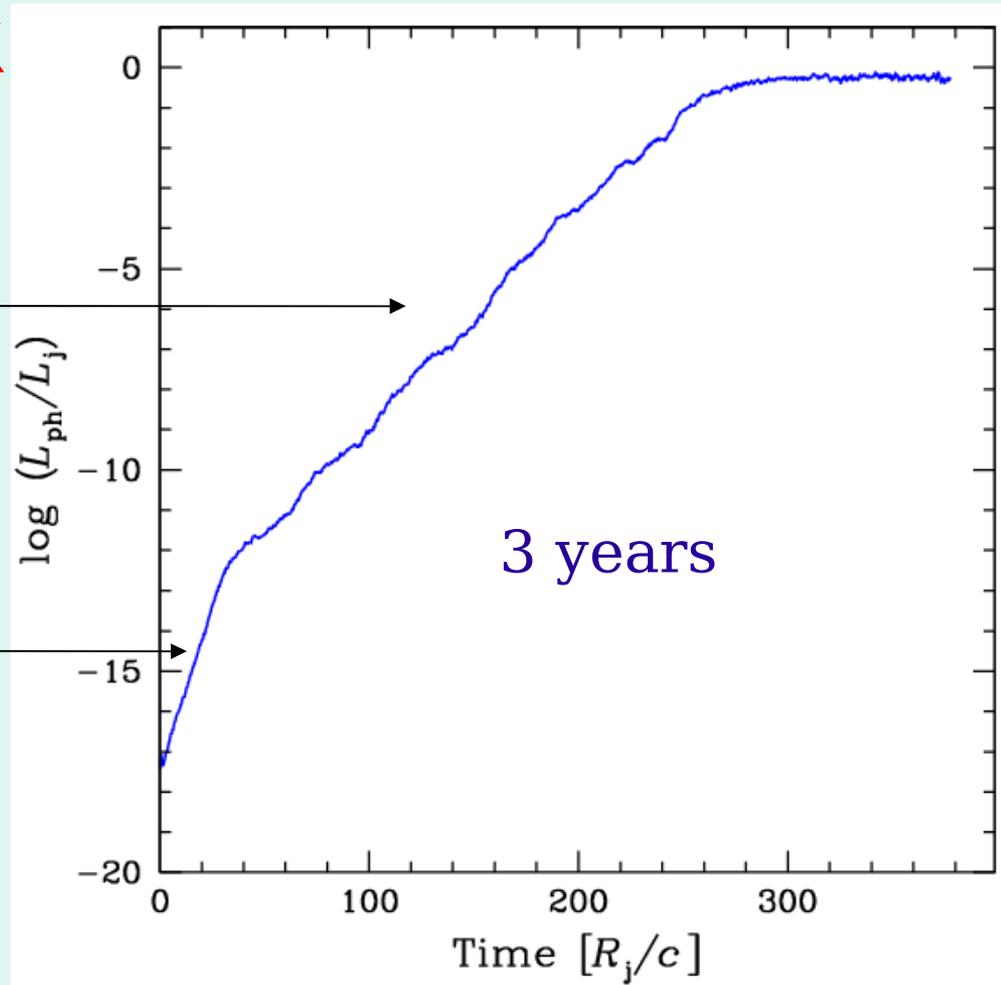
Exact 3D particle transport + space cells for the partner sampling

Exponential regime: Start of the breeding from extragalactic gamma-ray background



Breeding due to downstream-upstream feedback loop

Upstream - downstream photon avalanche



Results

$$R = 2 \cdot 10^{17} \text{ cm}$$

$\Gamma = 20$, magnetically

dominated $L_B = L_d$,

L	L_γ/L_j
$3 \cdot 10^{45} \text{ erg/s}$	0.25
10^{45}	0.32
$3 \cdot 10^{44}$	0.26
10^{44}	0.05
$3 \cdot 10^{43}$	0.

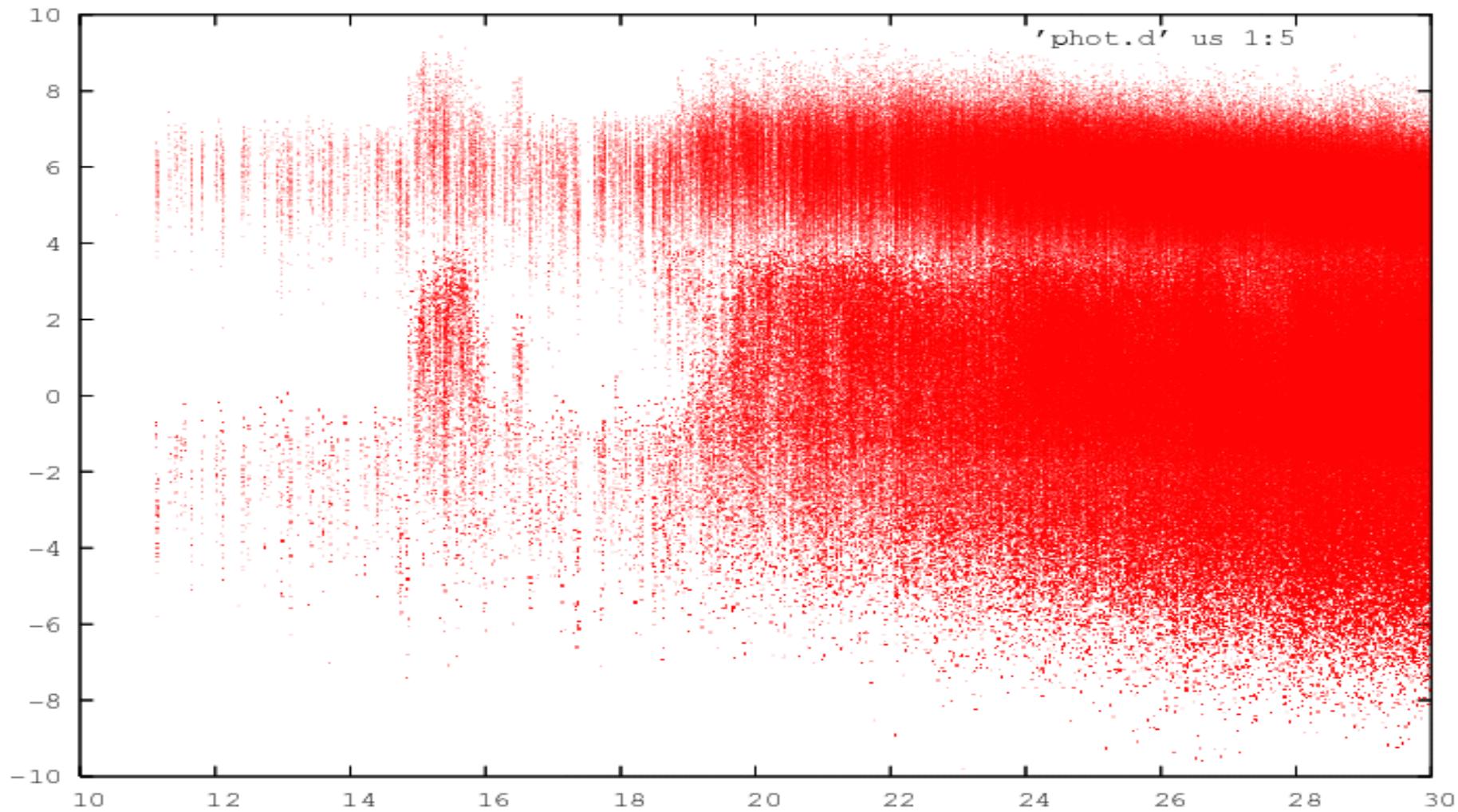
$\Gamma = 20$, $L_B = 0.2 L_d$

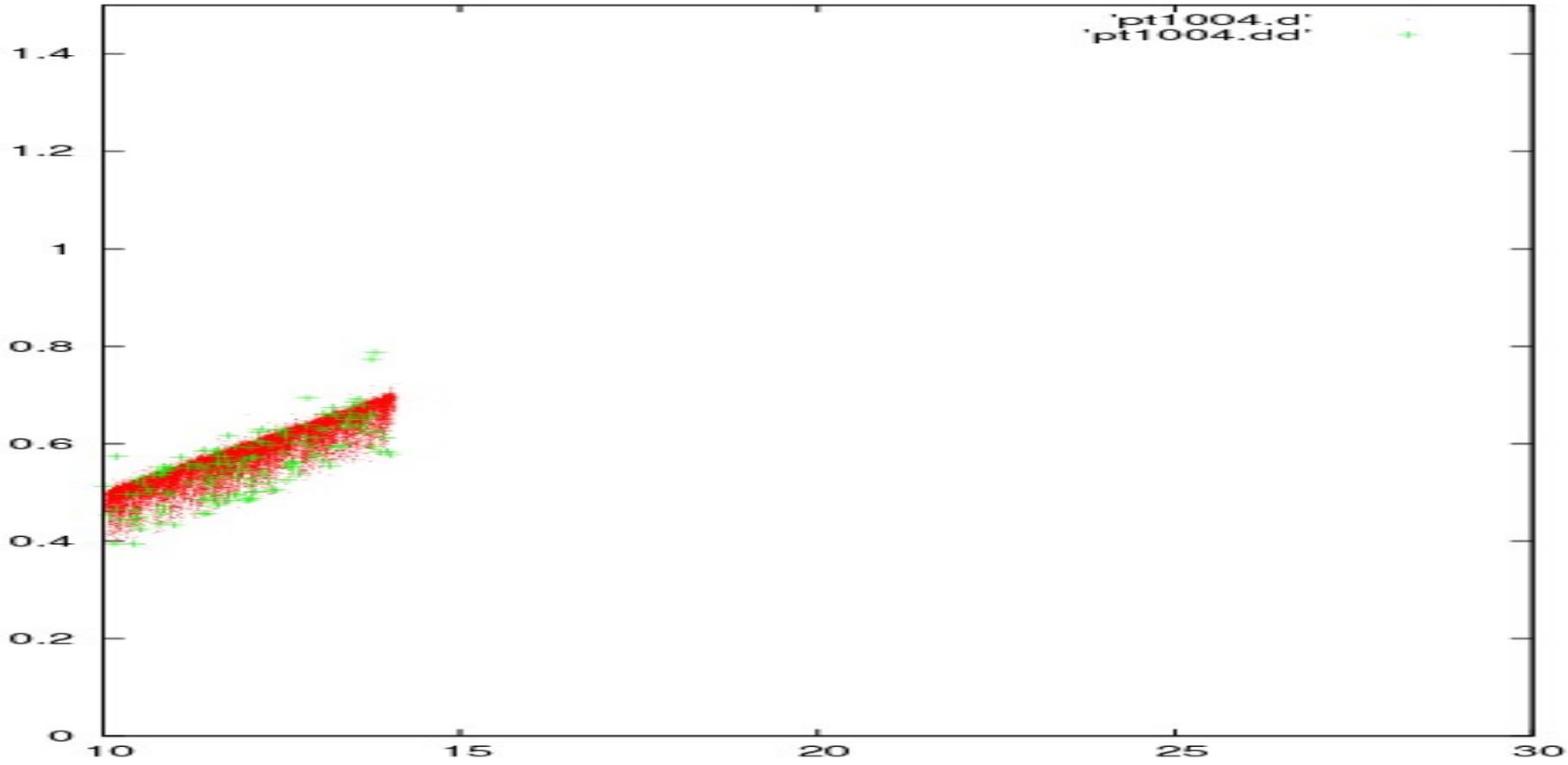
10^{46}	0.42
$3 \cdot 10^{45}$	0.44
10^{45}	0.50
$3 \cdot 10^{44}$	0.56
10^{44}	0.41
$5 \cdot 10^{43}$	0.12
$3 \cdot 10^{43}$	0

$L_B = 0.2 L_d$, $L_d = 3 \cdot 10^{44} \text{ erg/s}$

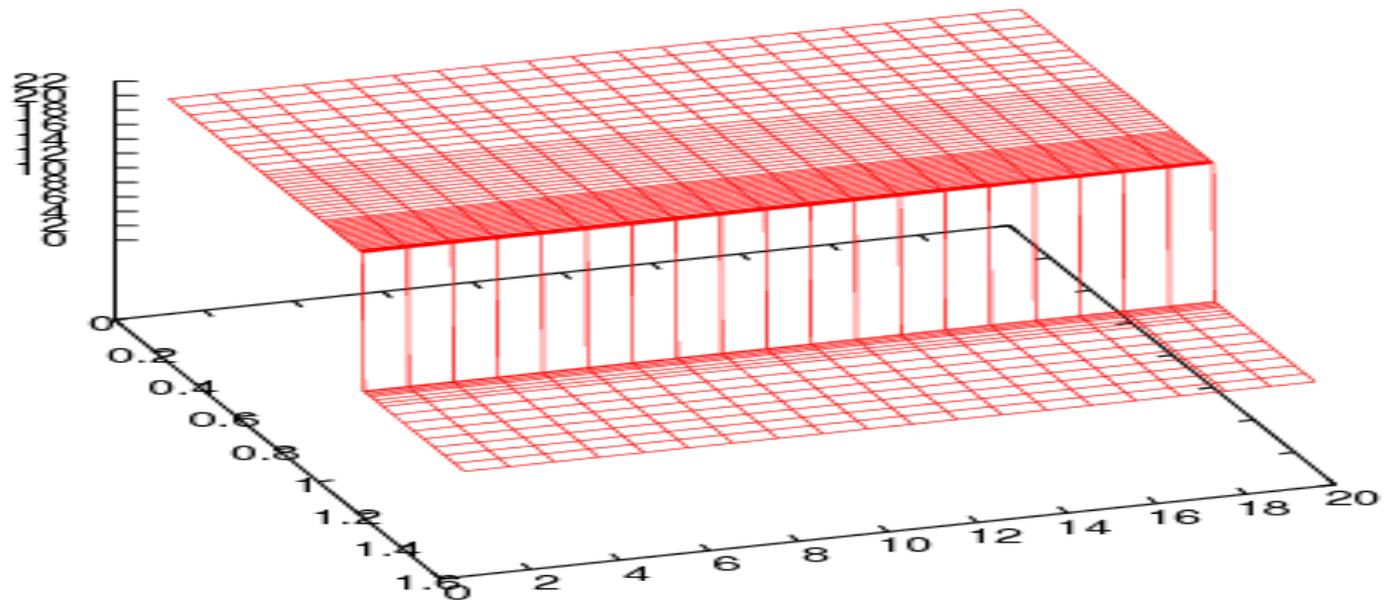
Γ	Γ_f	L_γ/L_j
40	11.5	0.81
30	12.4	0.76
20	13.6	0.56
14	13.6	0.12
12	12.	0.

Breeding works down
to $\Gamma = 8$ at a low
magnetization



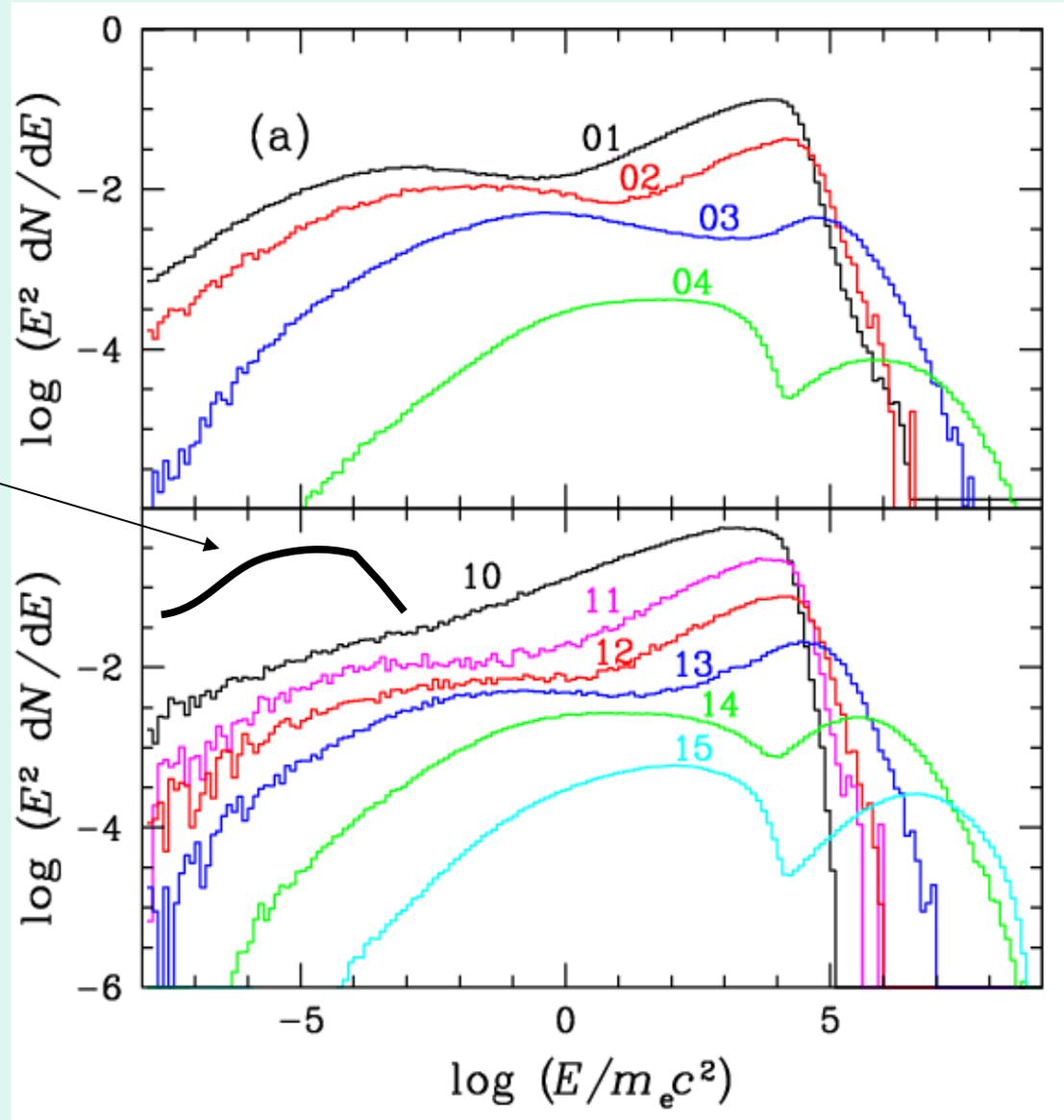


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Beam-on photon spectra “at infinity”.

An important detail is missing or poorly expressed: the “synchrotron peak”



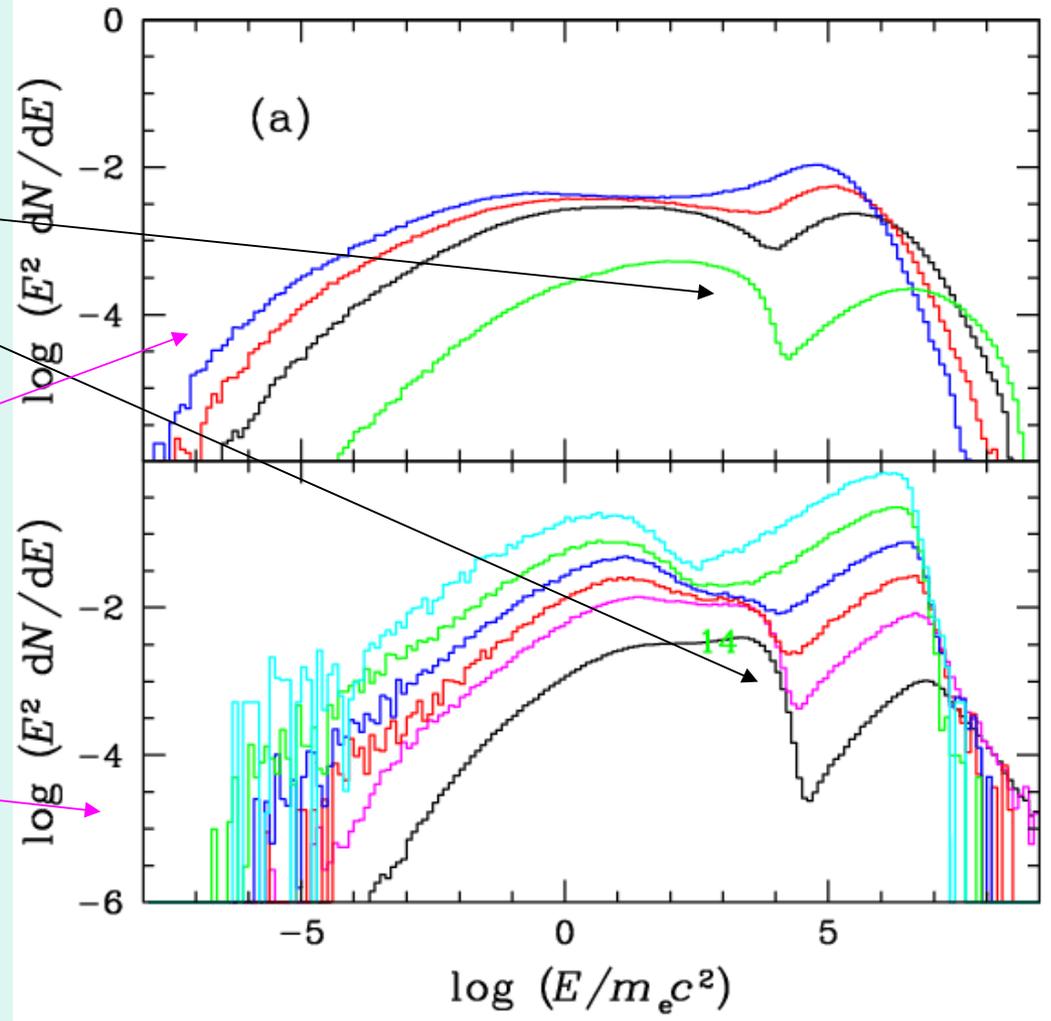
The ultimate synchrotron cutoff (Guilbert, Fabian, Rees, 1983)

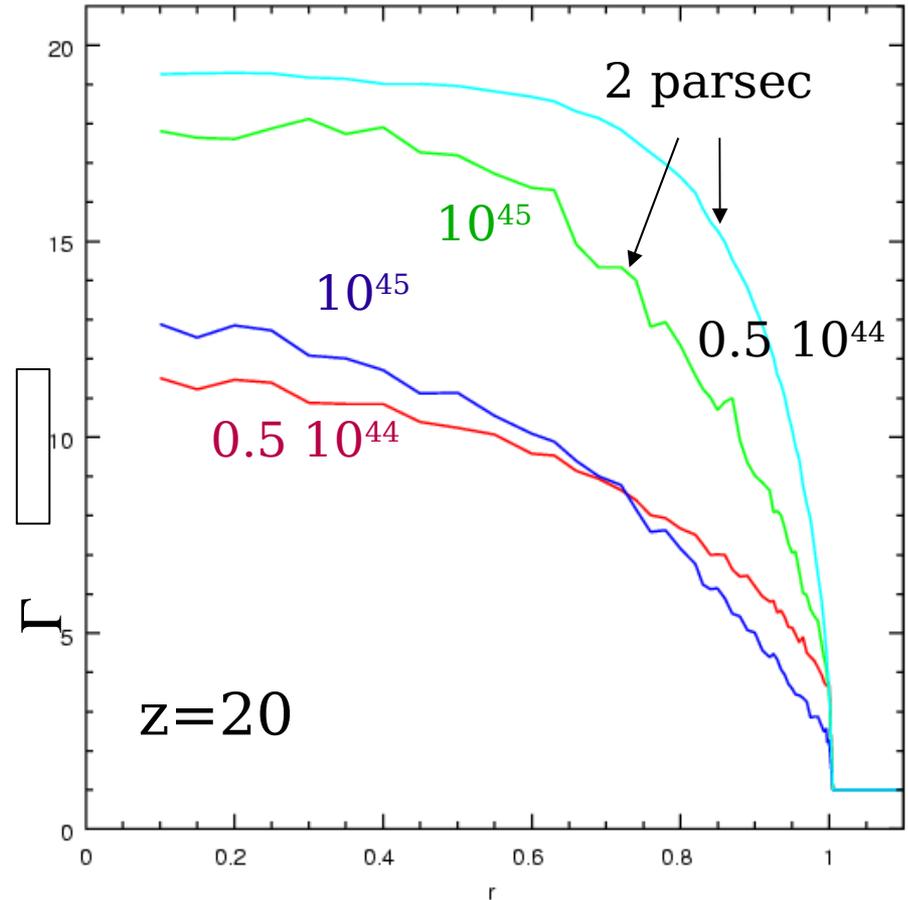
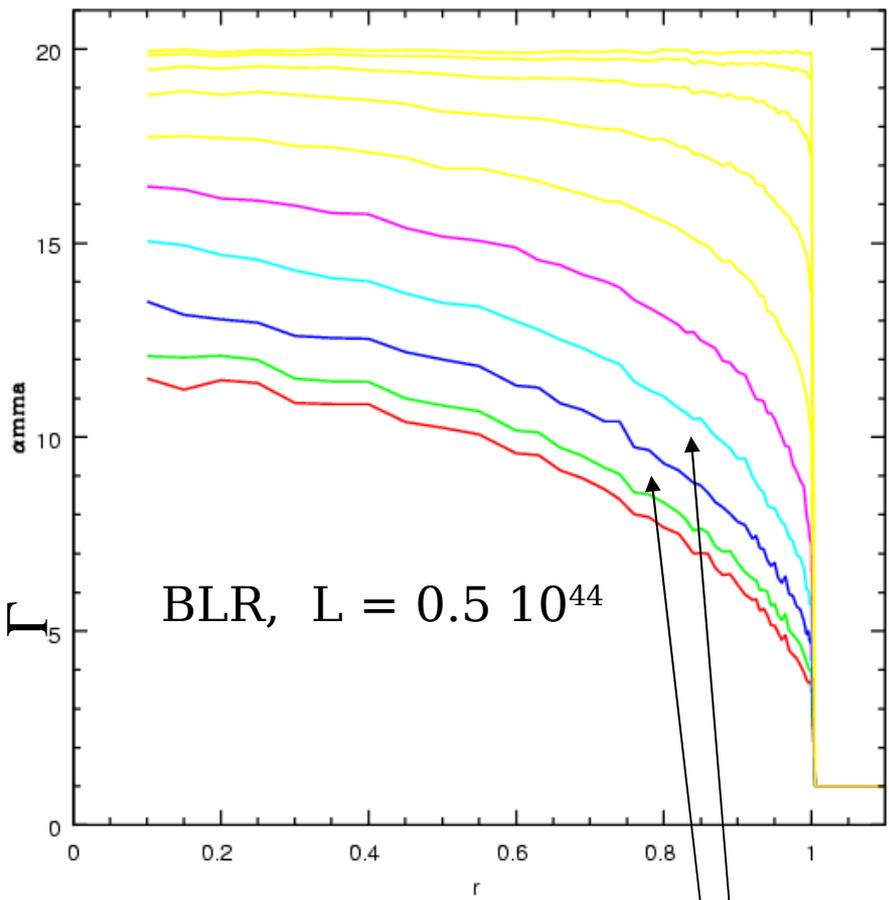
Low luminosity
spectra have a
feature observable by
Fermi:
a dip at several GeV

$$E_{s,\max} = m_e c^2 / \alpha \Gamma$$

$$F_t \sim 0.05 F_d$$

$$F_t \sim 0.01 F_d$$





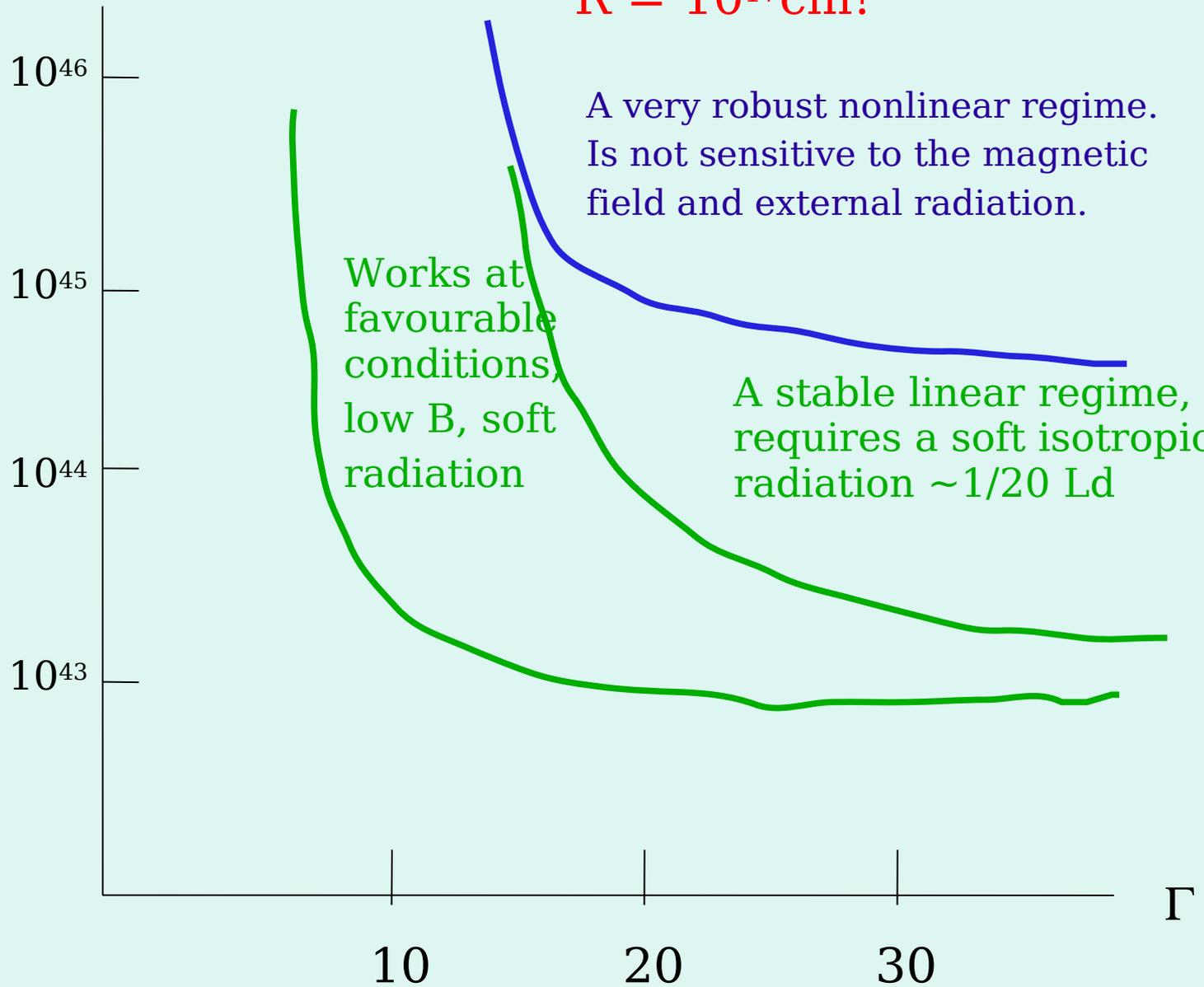
Radius

$\Delta z = 2$

Deceleration of the jet

L_d, L_j , erg/s

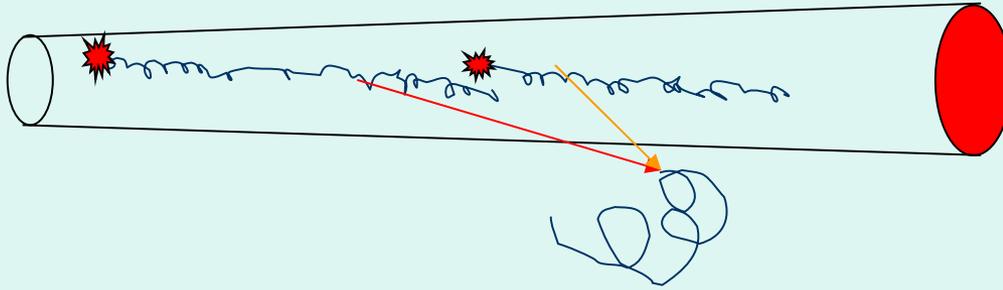
$R = 10^{17}$ cm!



- In the case of powerful jets of AGNs it is difficult to avoid the dissipation via the photon breeding.
- It is motivated by the requirement of a high efficiency of Gamma-ray emission.
- It, when taken alone, can not explain the entire spectrum of blazars. The photon breeding should be followed by a moderate pair reheating, which is quite natural.

What about BL Lacs?

Disk luminosity is insufficient to produce a large optical depth across the jet



Particle acceleration in the jet = input

Photon breeding = amplifier