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INTERACTIONS

Low-scale leptogenesis via neutrino oscillations

Inar Timiryasov

Niels Bohr Institute,
University of Copenhagen

International Conference on Particle Physics and Cosmology
dedicated to memory of Valery Rubakov

Yerevan, Armenia, 4th of October 2023



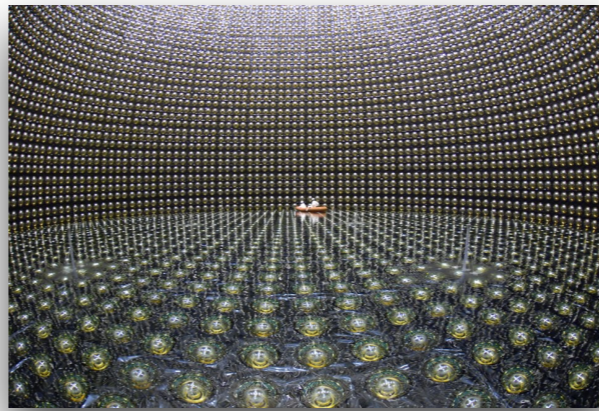
**At the first meeting of
the Chair of particle
physics and cosmology
at MSU**

13.11.2009

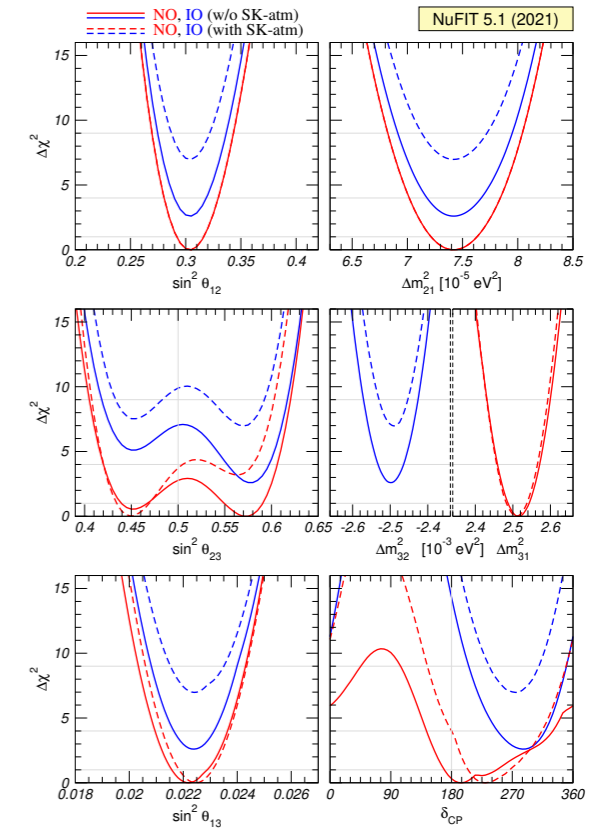
Beyond the Standard Model

- Neutrino flavour oscillations (violates L_α conservation, impossible if neutrinos are massless)

$$P(\nu_\alpha \rightarrow \nu_\beta) = \sin^2 2\theta \sin^2 \left(\frac{\Delta m^2 L}{4E} \right)$$

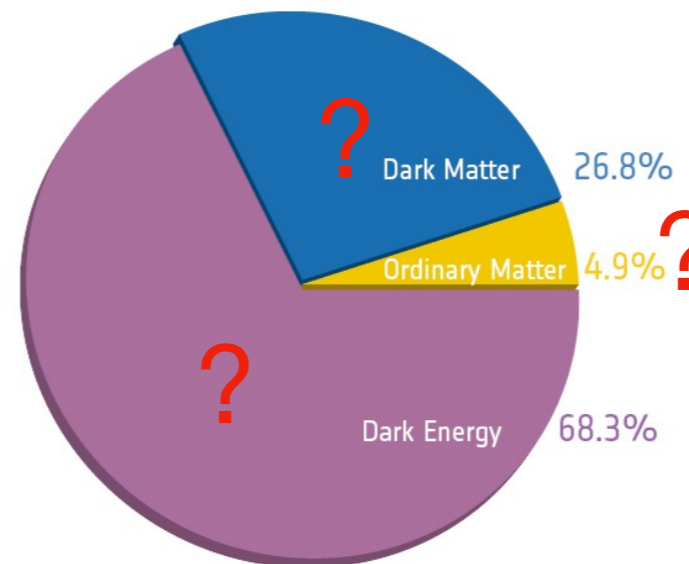


Super-Kamiokande (atmospheric oscillations $\nu_\mu \rightarrow \nu_\tau$)



NuFit collaboration <http://www.nu-fit.org>

- Cosmology



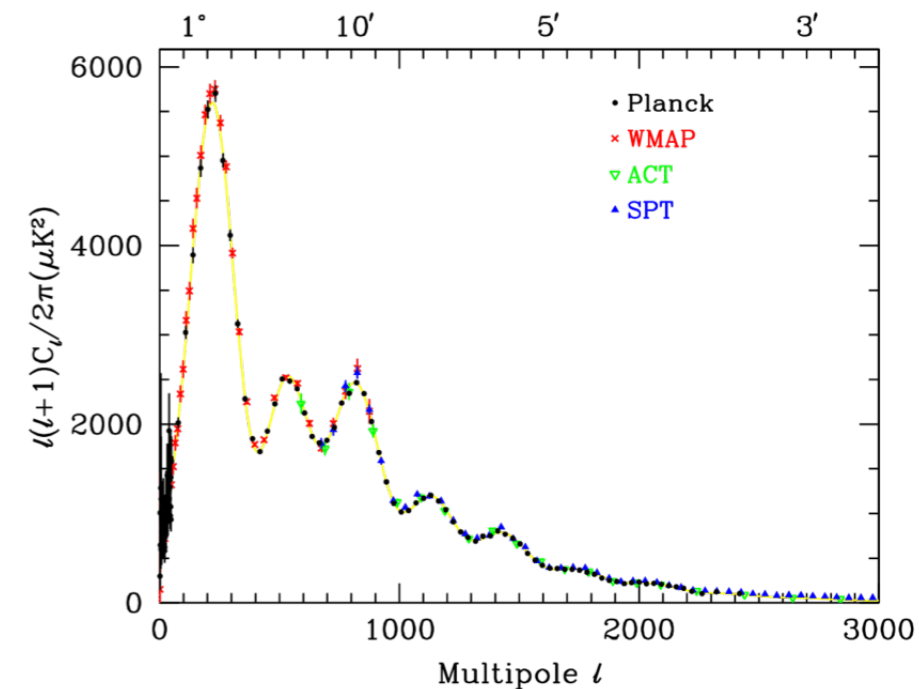
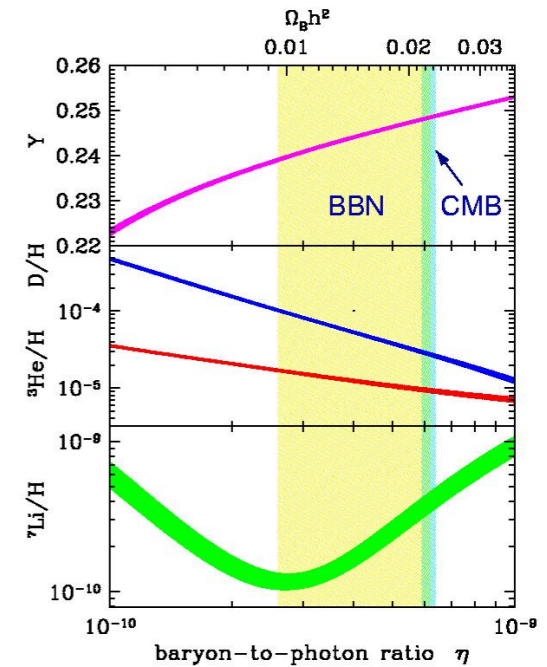
ESA and the Planck Collaboration

Baryon asymmetry of the Universe

- No antimatter in the present universe
- Baryon to photon ratio

$$\Delta = \frac{n_B - n_{\bar{B}}}{n_B + n_{\bar{B}}} \Bigg|_{T \sim 1 \text{ GeV}} \simeq \frac{n_B}{n_\gamma} \Bigg|_{\text{now}} \simeq 6 \times 10^{-10}$$

- At high T: $(10^{10} - 1)$ antiquarks per 10^{10} quarks
- Symmetric part annihilates into photons and ν
- Asymmetric part: origin of galaxies, stars, planets



Where the asymmetry comes from?

Sakharov Conditions (1967)

- Baryon number violation
- C and CP violation
- Deviation from thermal equilibrium

Where the asymmetry comes from?

Sakharov Conditions (1967)

- Baryon number violation

Nonperturbative sphaleron processes at $T > 130$ GeV
[Kuzmin, Rubakov, Shaposhnikov 1985]

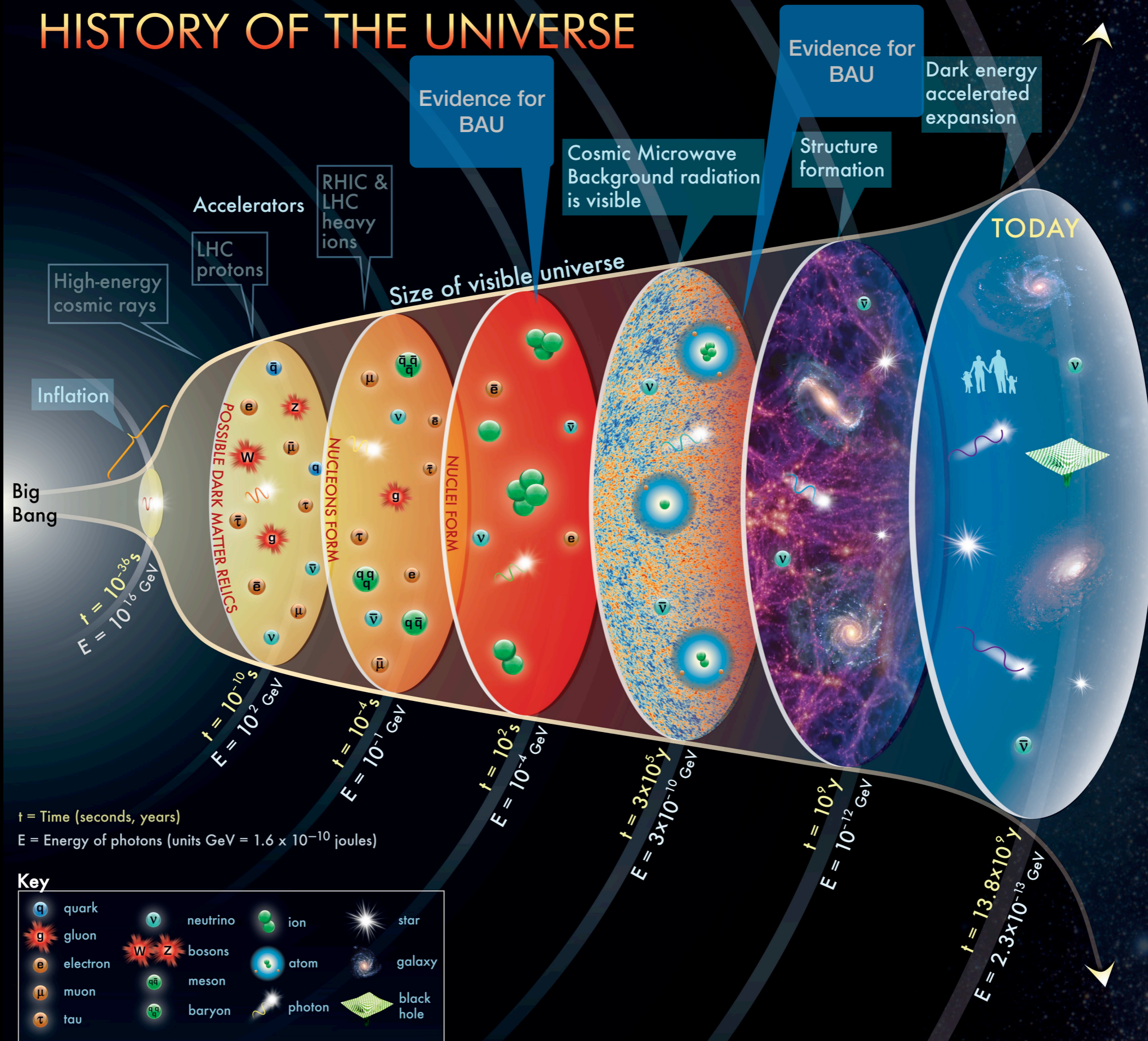
- C and CP violation

Present in the SM, but too small
 $G_F^6 s_1^2 s_2 s_3 \sin \delta m_t^4 m_b^4 m_c^2 m_s^2 \sim 10^{-20} \ll \Delta \sim 10^{-10}$

- Deviation from thermal equilibrium

No electroweak phase transition for $M_H > 73$ GeV
[Kajantie, Laine, Rummukainen, Shaposhnikov]

HISTORY OF THE UNIVERSE

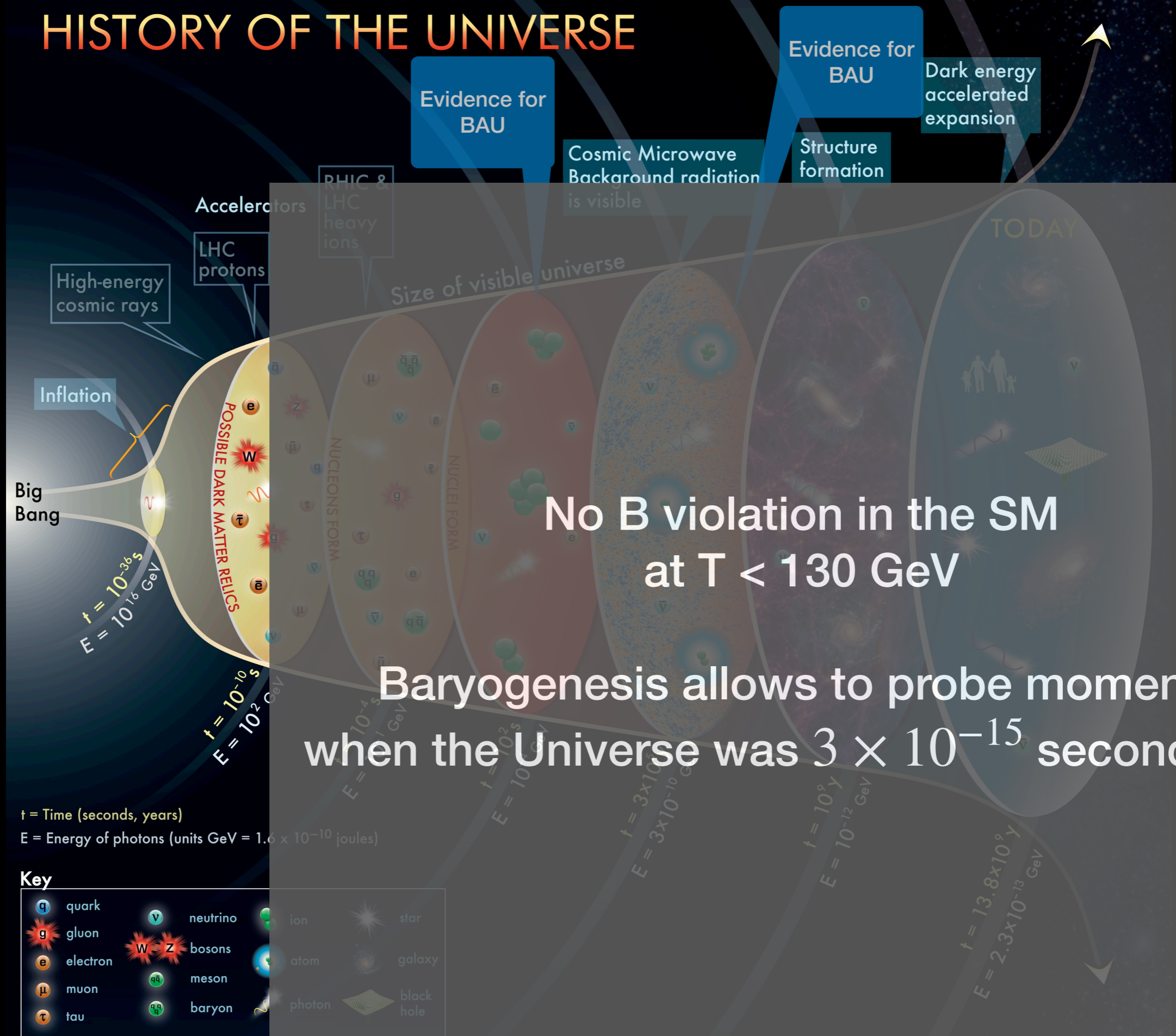


The concept for the above figure originated in a 1986 paper by Michael Turner.

Particle Data Group, LBNL © 2015

Supported by DOE

HISTORY OF THE UNIVERSE



Evidence for BAU

Evidence for BAU

Dark energy accelerated expansion

Cosmic Microwave Background radiation is visible

Structure formation

Accelerators

LHC protons

RHIC & LHC heavy ions

High-energy cosmic rays

Inflation

Big Bang

$t = 10^{-36} s$
 $E = 10^{16} GeV$

$t = 10^{-10} s$
 $E = 10^2 GeV$

Size of visible universe

No B violation in the SM
at $T < 130 GeV$

Baryogenesis allows to probe moments
when the Universe was 3×10^{-15} seconds old

t = Time (seconds, years)

E = Energy of photons (units GeV = 1.6×10^{-10} joules)

Key

	quark		neutrino		ion		star
	gluon		bosons		atom		galaxy
	electron		meson		photon		black hole
	muon		baryon				
	tau						

The concept for the above figure originated in a 1986 paper by Michael Turner.

Baryogenesis via neutrino oscillations

E. Kh. Akhmedov^(a,b) V. A. Rubakov^(c,a,d) and A. Yu. Smirnov^(a,c)

^(a)The Abdus Salam International Centre for Theoretical Physics, I-34100 Trieste, Italy

^(b)National Research Centre Kurchatov Institute, Moscow 123182, Russia

^(c)Institute for Nuclear Research of the Russian Academy of Sciences, Moscow 117312, Russia

^(d)Institute for Cosmic Ray Research, University of Tokyo, Tanashi, Tokyo 188, Japan

(March 5, 1998)

We propose a new mechanism of leptogenesis in which the asymmetries in lepton numbers are produced through the CP-violating oscillations of “sterile” (electroweak singlet) neutrinos. The asymmetry is communicated from singlet neutrinos to ordinary leptons through their Yukawa couplings. The lepton asymmetry is then reprocessed into baryon asymmetry by electroweak sphalerons. We show that the observed value of baryon asymmetry can be generated in this way, and the masses of ordinary neutrinos induced by the seesaw mechanism are in the astrophysically and cosmologically interesting range. Except for singlet neutrinos, no physics beyond the Standard Model is required.

PACS: 98.80.Cq, 14.60.St

IC/98/22, INR-98-14T

[hep-ph/9803255](https://arxiv.org/abs/hep-ph/9803255)

<https://arxiv.org/abs/hep-ph/9803255>

The seesaw mechanism

$$\mathcal{L} = \mathcal{L}_{SM} + i\bar{\nu}_{R_I}\gamma^\mu\partial_\mu\nu_{R_I} - F_{\alpha I}\bar{L}_\alpha\tilde{\Phi}\nu_{R_I} - \frac{M_{IJ}}{2}\bar{\nu}_{R_I}^c\nu_{R_J} + h.c.$$

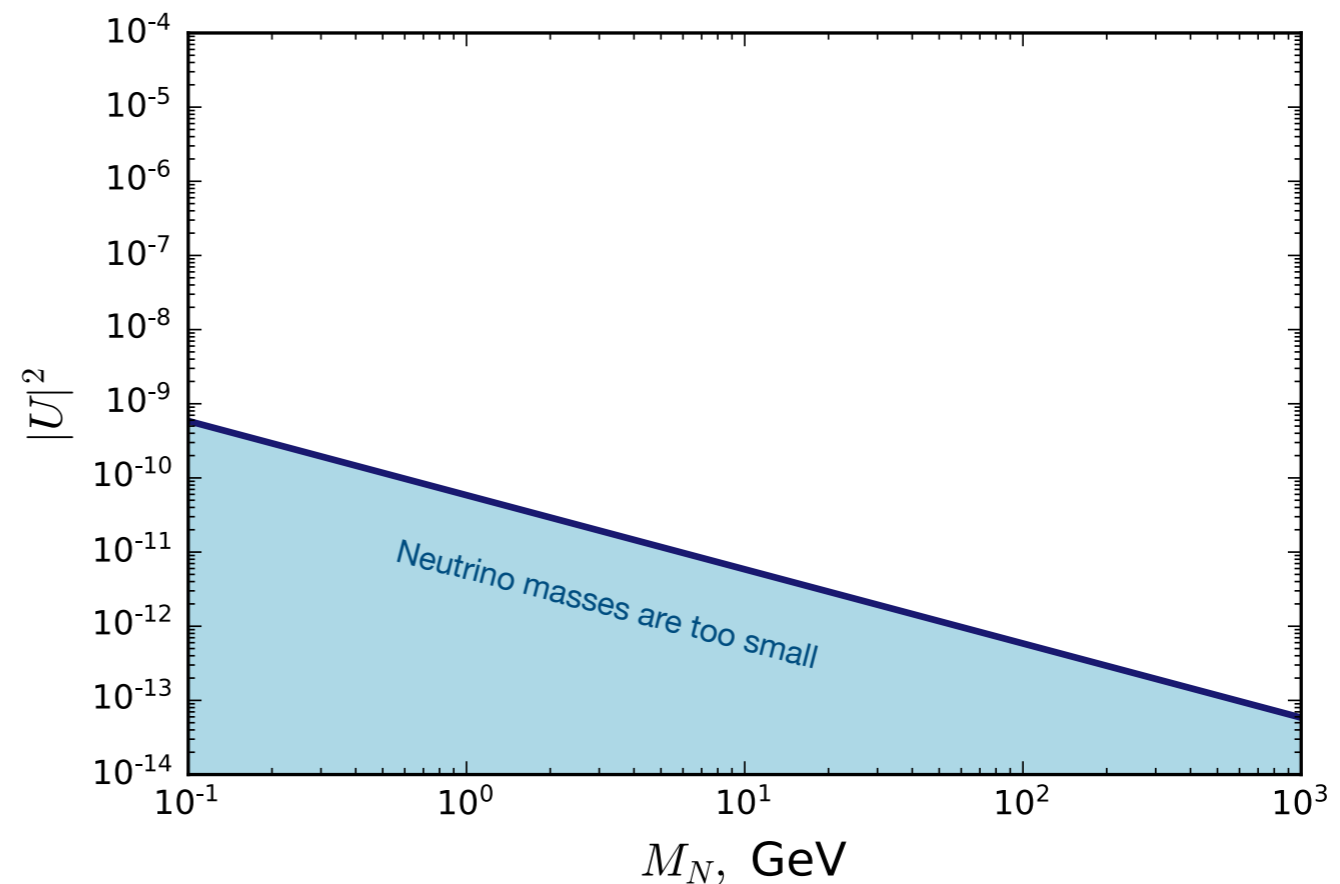
Minkowski; Yanagida; Gell-Mann,
Ramond, Slansky; Glashow;
Mohapatra, Senjanovic

eV sterile neutrinos are
outside of this range

The plenary talk by
Mikhail Danilov

- New singlet fermions
- Mixing with light neutrinos

$$\nu_{L_\alpha} = U_{\alpha i}^{PMNS}\nu_i + \Theta_{\alpha I}N_I^c$$



We consider nearly degenerate HNLs (Heavy Neutral Leptons)

Heavy Neutral Leptons: Leptogenesis

N can be responsible for the Baryon Asymmetry

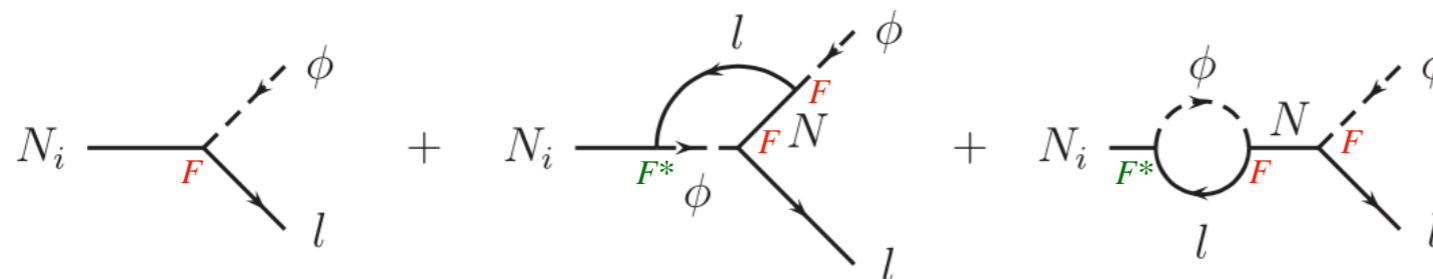
Fukugita and Yanagida, 1986

Reviews: Buchmuller, Di Bari, Plumacher:

Leptogenesis for pedestrians, 2004

Bödeker, Buchmuller, 2009.07294

- B violated by sphaleron processes
- CP asymmetry in N decays
- Deviation from equilibrium when $\Gamma_N \sim H$



$$\varepsilon_i = \frac{\Gamma(N_i \rightarrow l\phi) - \Gamma(N_i \rightarrow \bar{l}\bar{\phi})}{\Gamma(N_i \rightarrow l\phi) + \Gamma(N_i \rightarrow \bar{l}\bar{\phi})}$$

$$\varepsilon \sim \frac{\text{Im}(F^\dagger F)^2}{|F|^2}$$

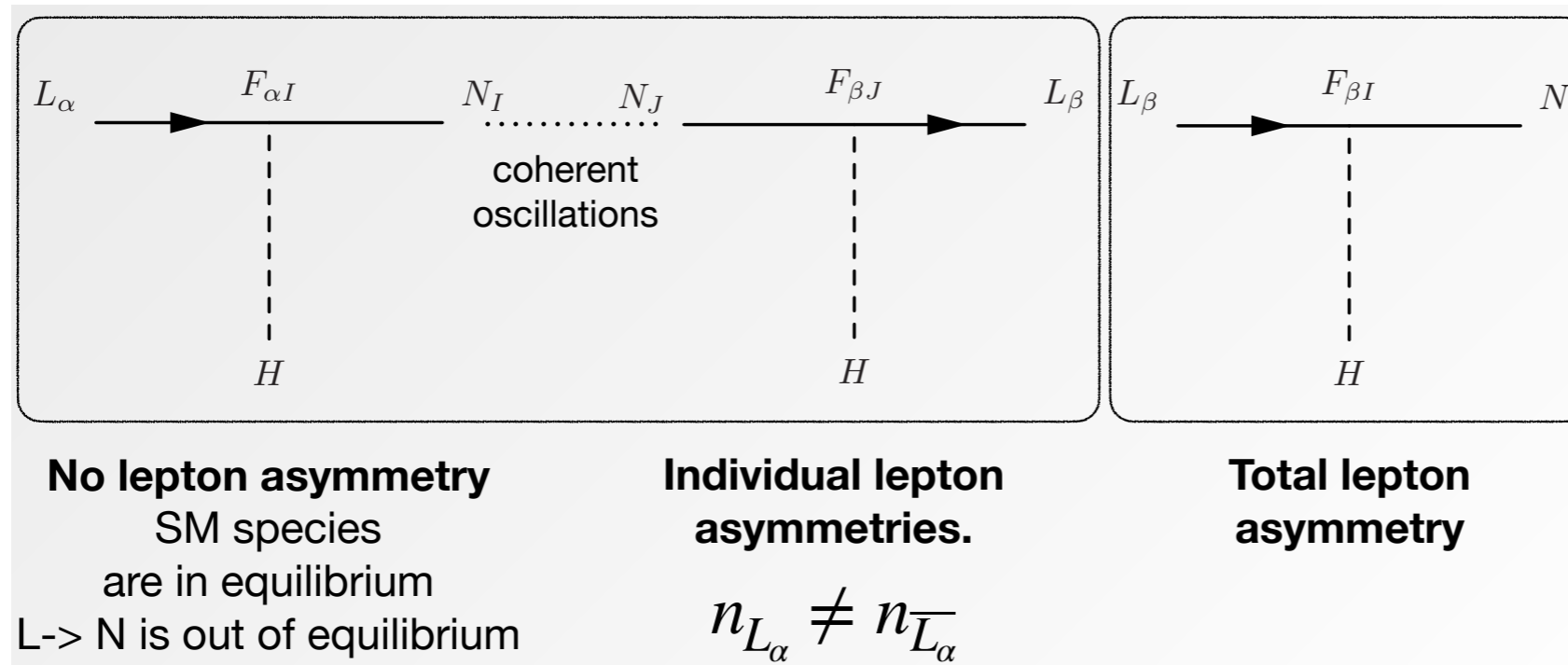
Davidson Ibarra bound, 2002

$$M \gtrsim 10^9 \text{ GeV}$$

$$\varepsilon_{\text{max}} = \frac{3}{16\pi} \frac{M m_{\text{atm}}}{v^2} \simeq 10^{-6} \left(\frac{M}{10^{10} \text{ GeV}} \right)$$

Low-scale leptogenesis via neutrino oscillations

- B violated by sphaleron processes
- **CP asymmetry is enhanced by N-N oscillations**
- Deviation from equilibrium: small Yukawas
(masses are also relatively small — “low-scale”)



Akhmedov, Rubakov, Smirnov 1998

Asaka, Shaposhnikov 2005

Canetti, Drewes, Frossard; Eijima, Ishida; Shuve, Yavin; Abada, Arcadi, Domcke, Lucente; Hernández, Kekic, López-Pavón, Racker, Salvado; Drewes, Garbrecht, Gueter, Klaric; Hambye, Teresi; Ghiglieri, Laine; IT; ...

Description of low-scale leptogenesis

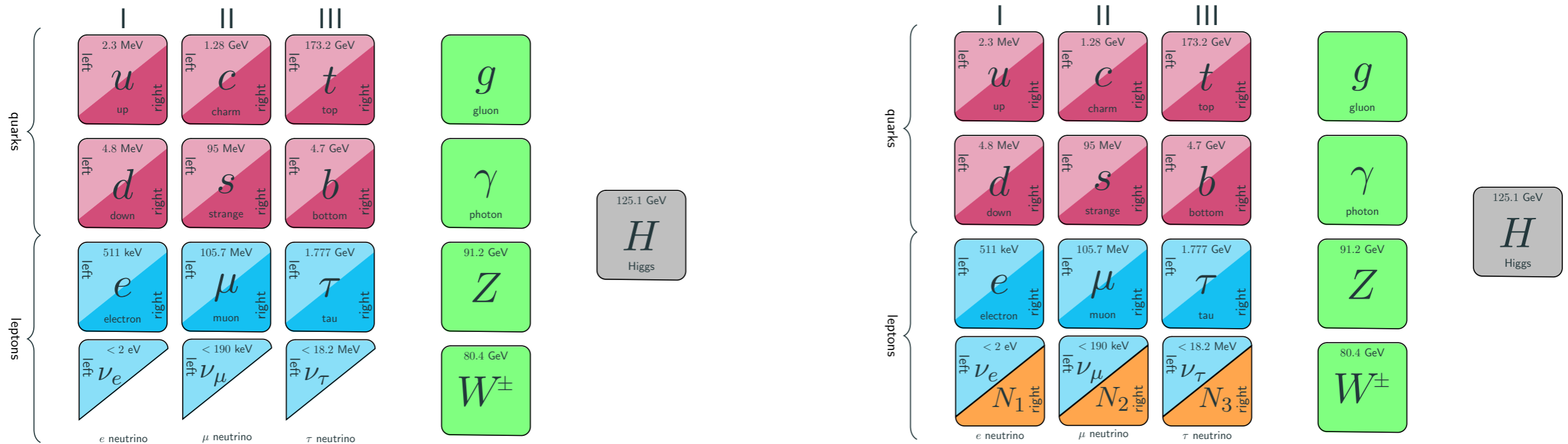
Significant theoretical developments since 2014

[1605.07720, 1703.06085, 1703.06087, 1605.07720, 1709.07834, 1711.08469, 1208.4607, 1606.06690, 1606.06719, 1609.09069, 1710.03744, 1808.10833, 1811.01971, 1905.08814, 1911.05092, 2004.10766, 2008.13771, 2203.05772]

- Fermion number violating processes (processes with and without helicity flip)
Eijima, Shaposhnikov; Ghiglieri, Laine
- Accurate computation of the rates (including Landau-Pomeranchuk-Migdal resummation of multiple soft scatterings)
Ghiglieri, Laine
- Spectator processes
Shuve, Yavin; Ghiglieri, Laine; Eijima, Shaposhnikov, IT
- Gradual sphaleron freeze-out
Ghiglieri, Laine; Eijima, Shaposhnikov, IT
- Rates for HNLs with $M \sim M_W$
Klaric, Shaposhnikov, IT

Neutrino Minimal Standard Model (ν MSM)

Asaka, Blanchet, Shaposhnikov 2005
Asaka, Shaposhnikov 2005



N_1

DM candidate

Monday talk by Misha Shaposhnikov

N_2

ν masses via see-saw
BAU
(DM production)

N_3

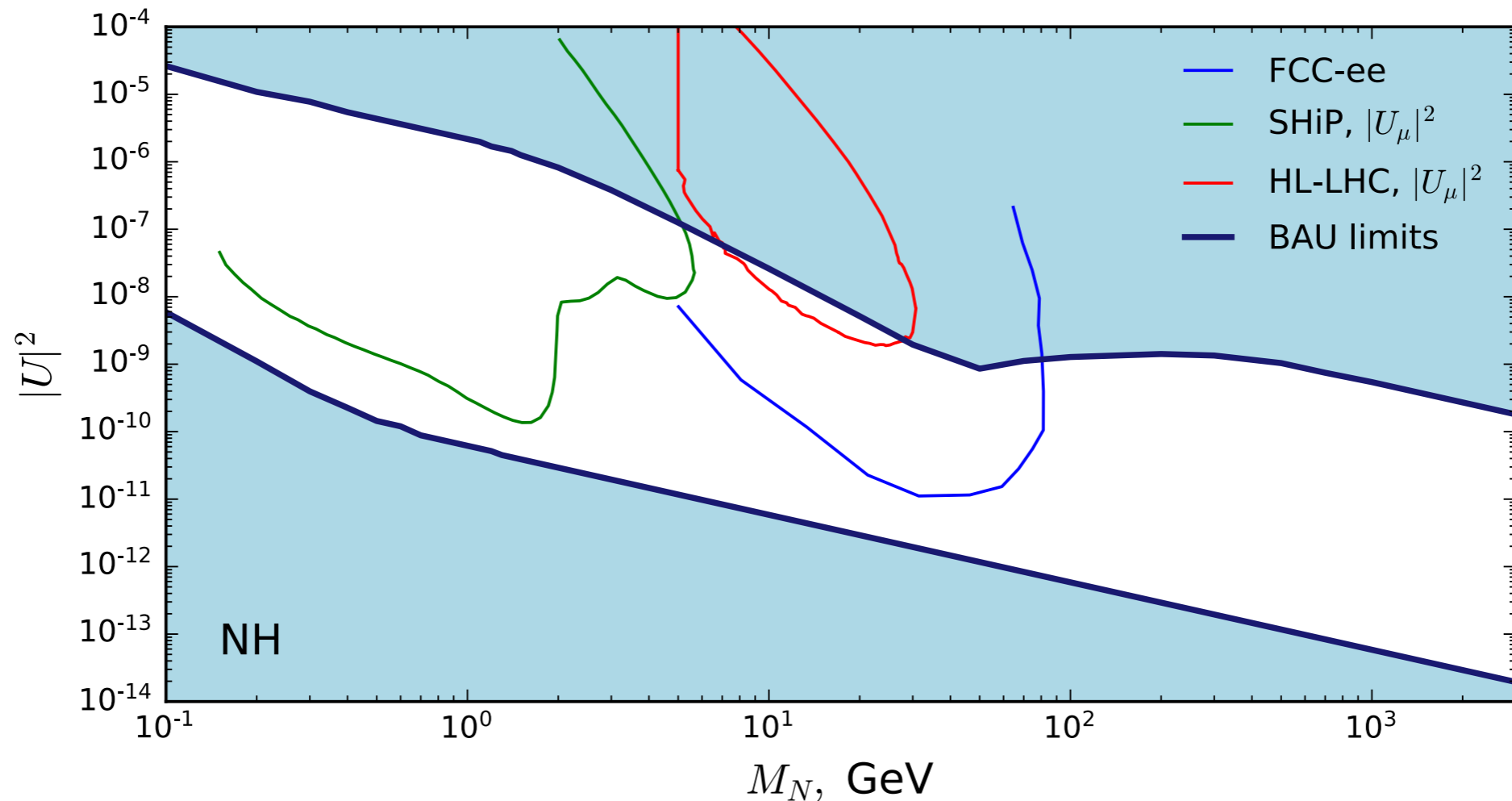
$$M_N \gtrsim 0.1 \text{ GeV}$$

Nearly degenerate

Baryogenesis via oscillations
Akhmedov, Rubakov, Smirnov, 1998
Asaka, Shaposhnikov 2005

Uniting leptogeneses

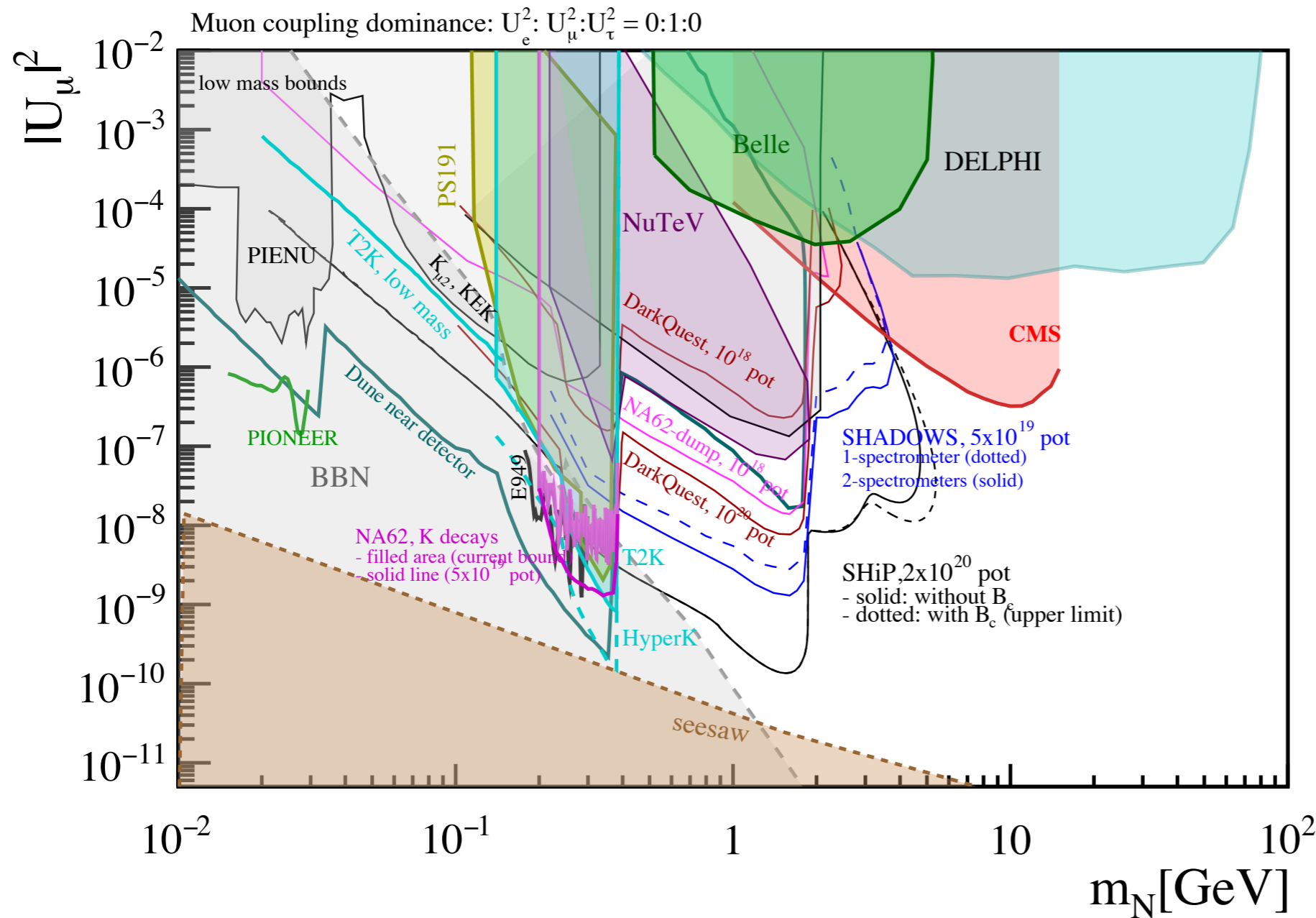
Juraj Klarić, Mikhail Shaposhnikov, IT [2008.13771](#), Phys.Rev.Lett. 127 (2021)



- Leptogenesis via oscillations still works for heavy HNLs because the washout of the asymmetry can vary a lot for different lepton flavours (*flavour hierarchical washout*)
- Resonant leptogenesis works for $M_N \gtrsim 5$ GeV since the asymmetry generated in HNL decays into a certain flavour can be very large

3RH case: Klarić, Georis, Drewes [2106.16226](#)

The quest for Heavy Neutral Leptons

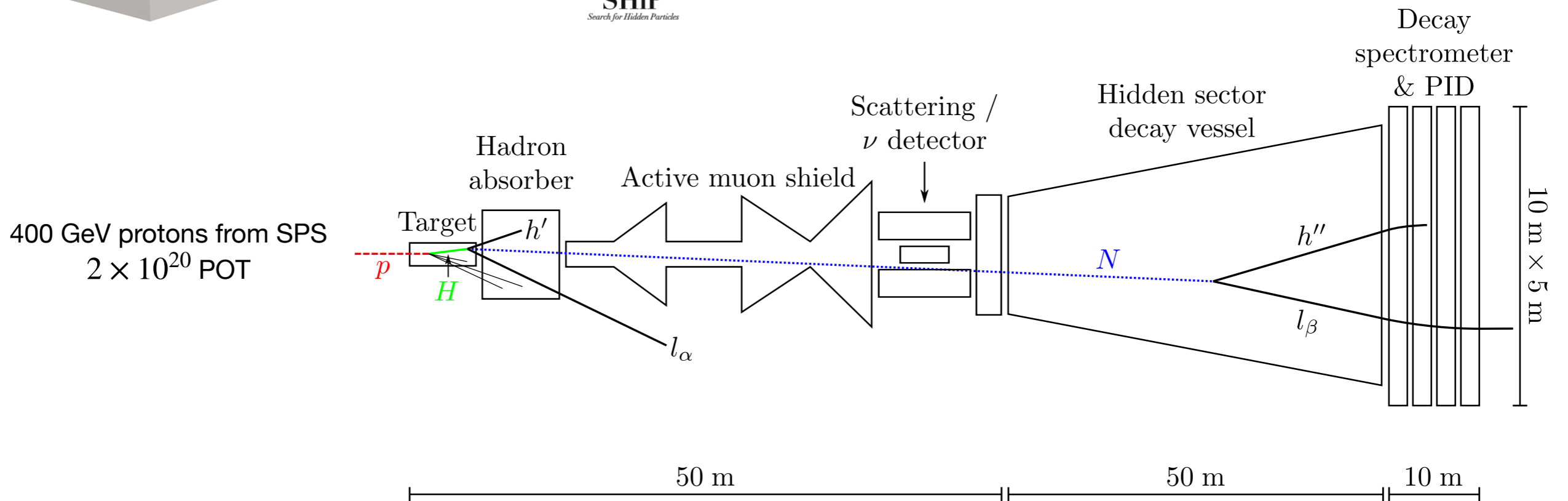
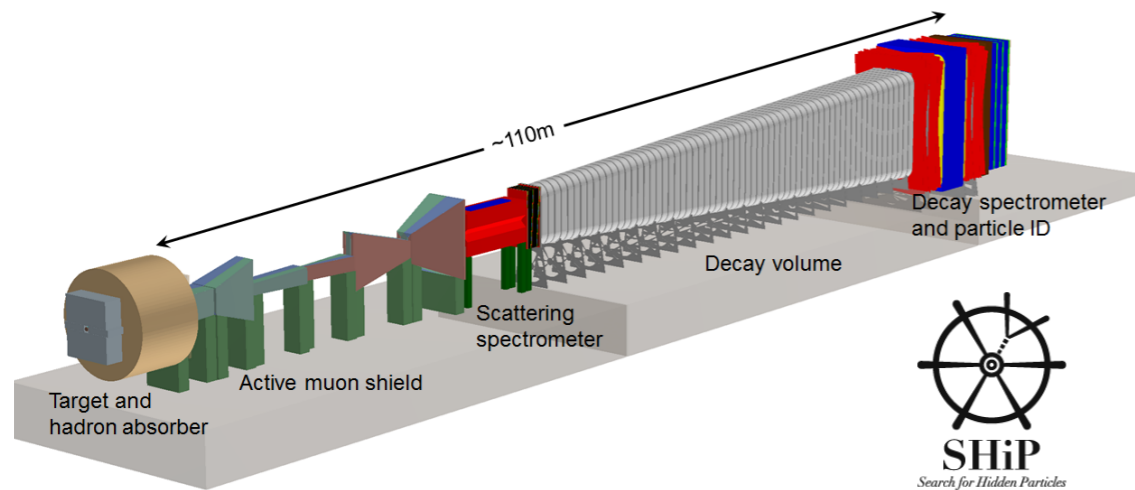


The Present and Future Status of Heavy Neutral Leptons
2203.08039

For a unified sensitivity estimation:
a new Mathematica package SensCalc
<https://arxiv.org/abs/2305.13383>
<https://doi.org/10.5281/zenodo.7957784>

How to search for HNLs?

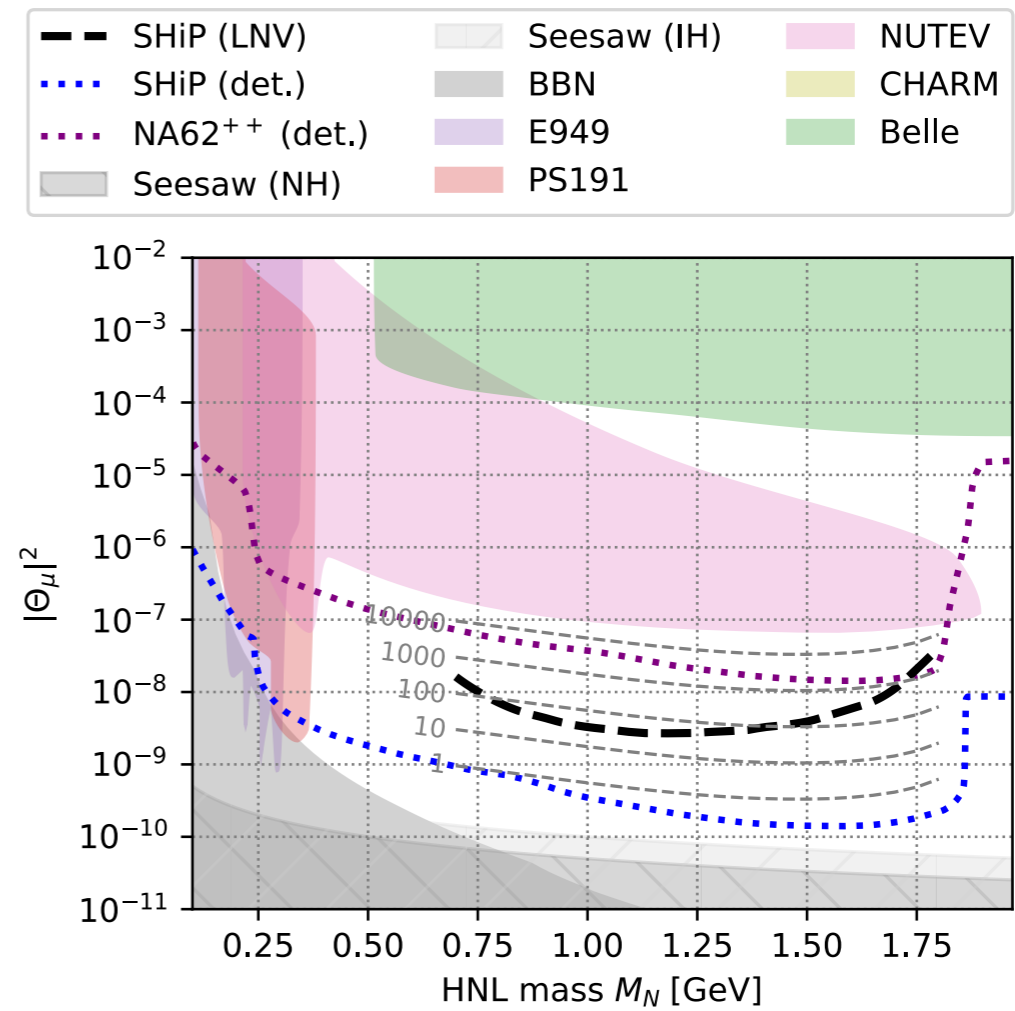
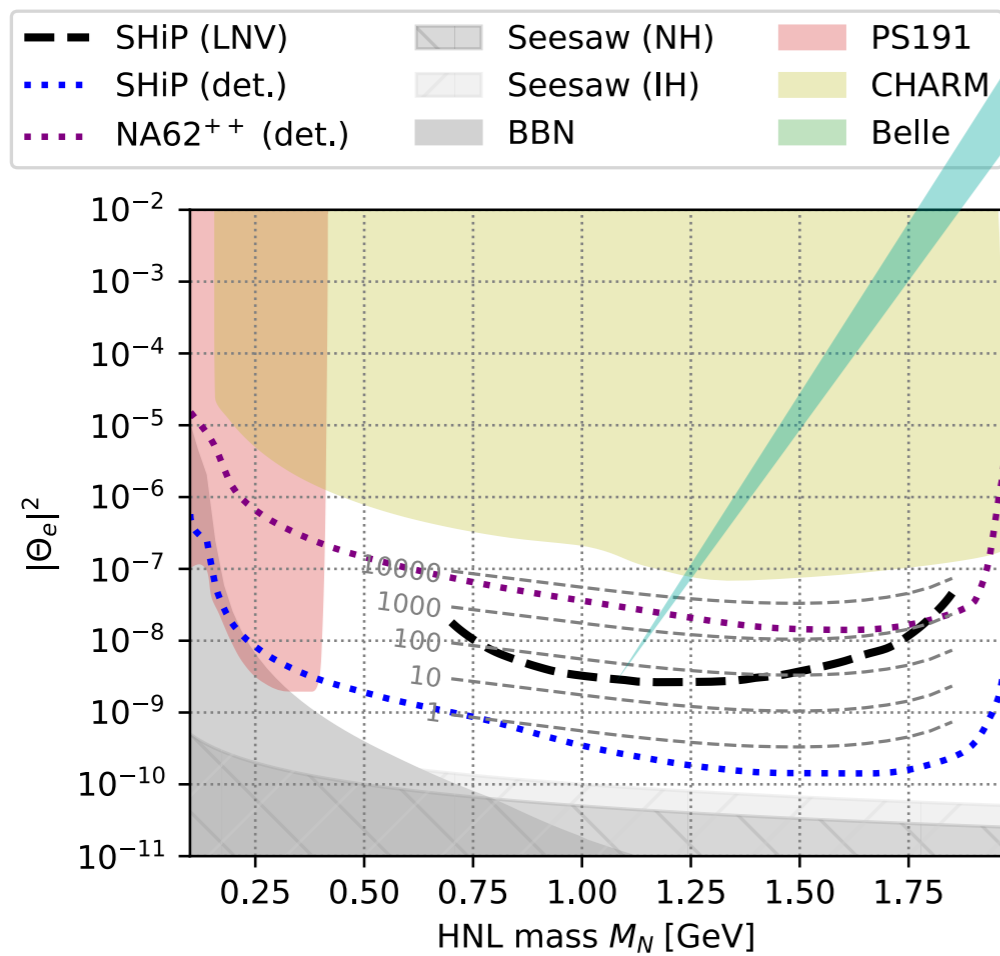
Example: SHiP - Search for Hidden Particles



* I am a member of SHiP collaboration

Probing lepton number violation at SHiP

Above this line SHiP can distinguish Majorana vs Dirac



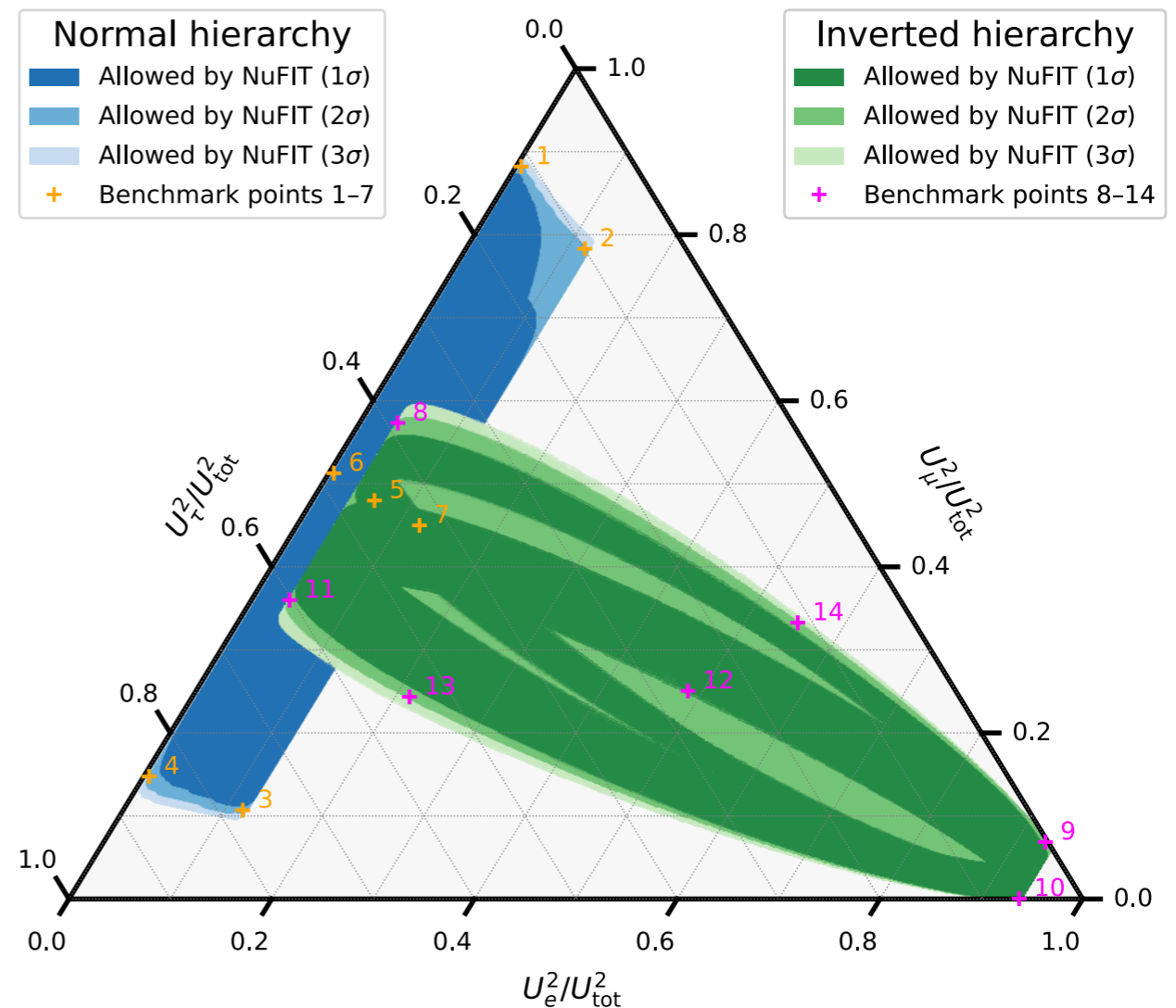
Neutrino oscillation data and mixings

Many analyses assume mixing with a single neutrino

Not all mixing angles are allowed in the model with two HNLs

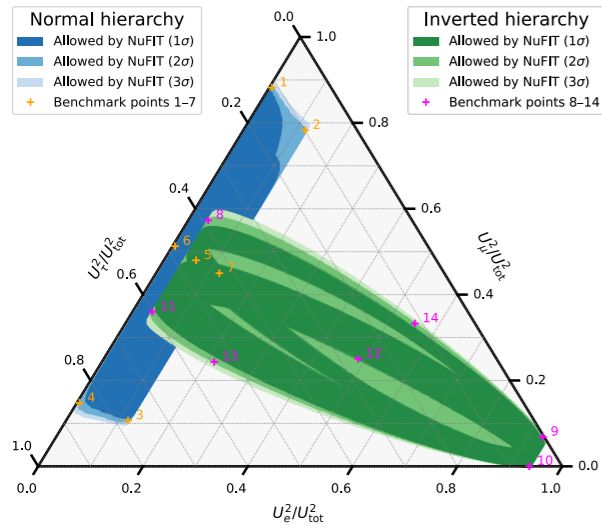
$$\nu_{L\alpha} = U_{\alpha i}^{PMNS} \nu_i + \Theta_{\alpha I} N_I^c$$

$$U_\alpha^2 \equiv \sum_I |\Theta_{\alpha I}|^2 \quad \text{and} \quad U_{\text{tot}}^2 \equiv \sum_{\alpha, I} |\Theta_{\alpha I}|^2$$

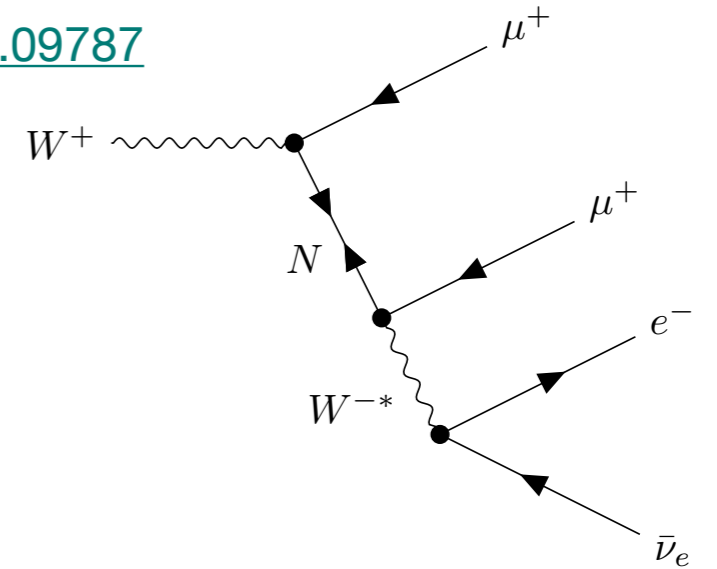


$$U_e^2/U_{\text{tot}}^2 + U_\mu^2/U_{\text{tot}}^2 + U_\tau^2/U_{\text{tot}}^2 = 1$$

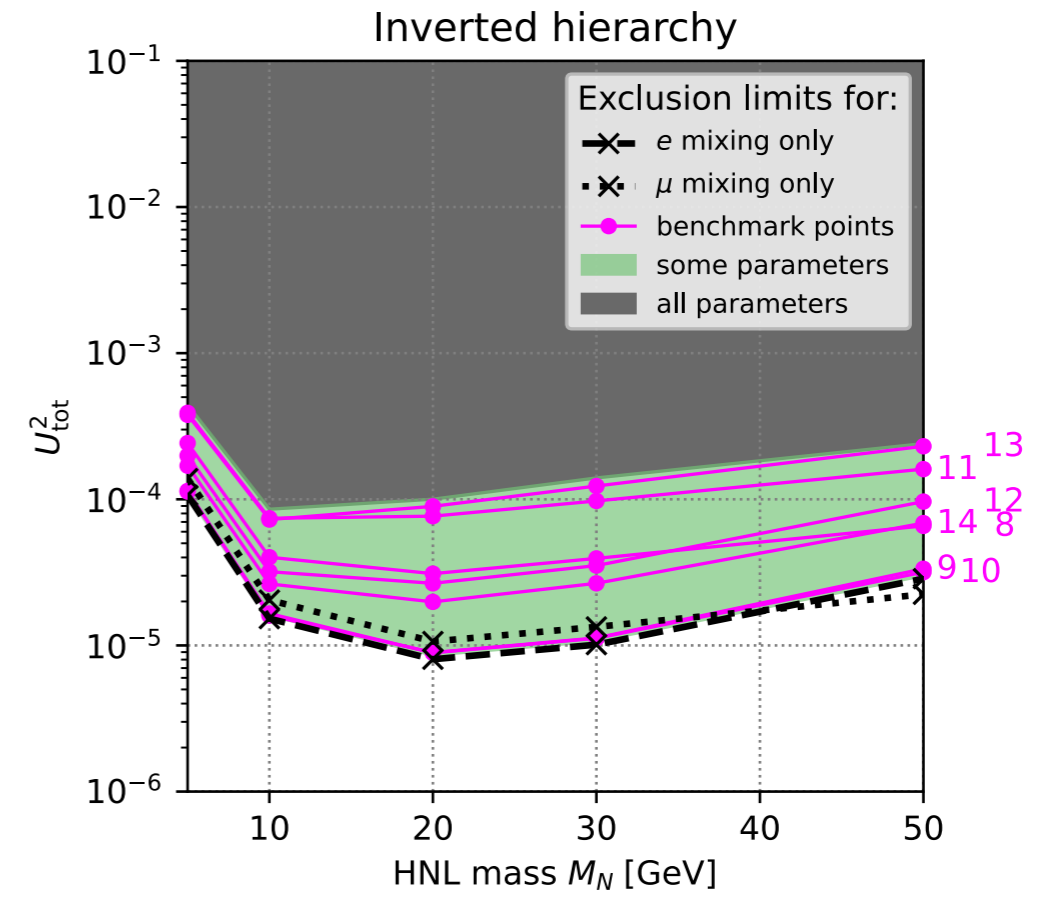
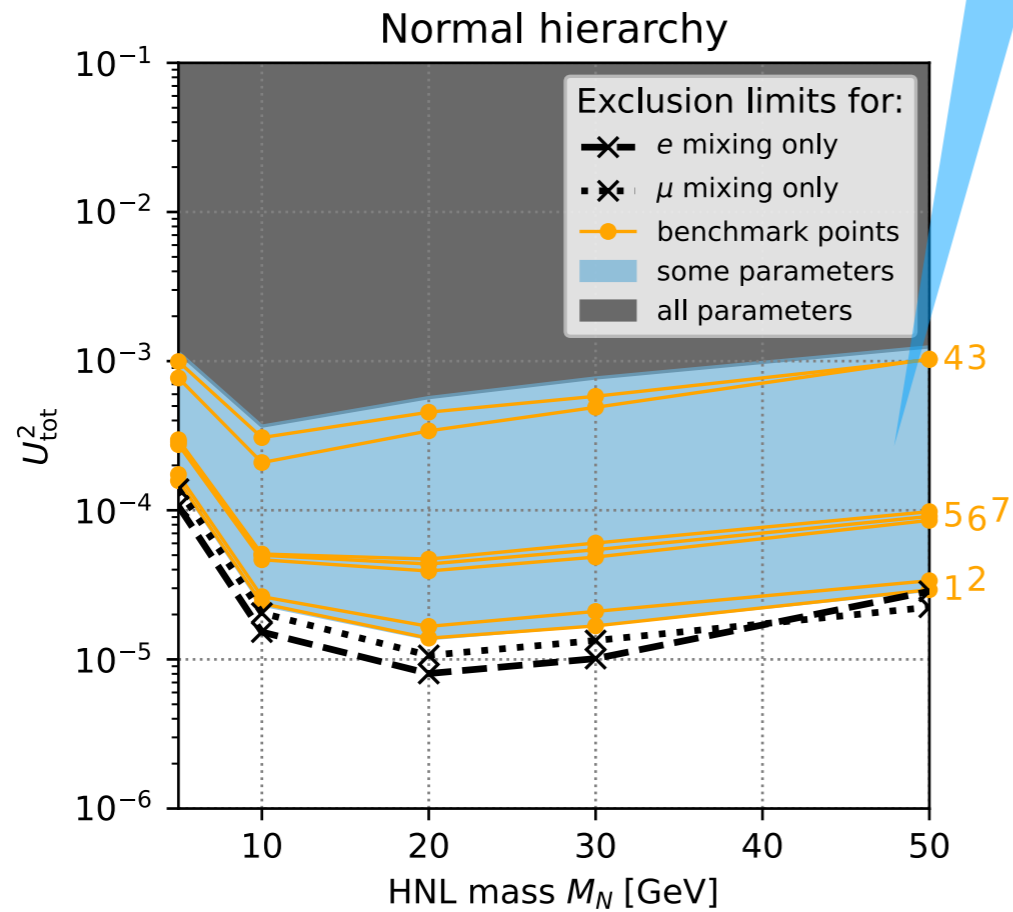
Neutrino oscillation data and reinterpretation



ATLAS triplepton search [1905.09787](https://arxiv.org/abs/1905.09787)



New physics could hide here!



Jean-Loup Tastet, Oleg Ruchayskiy, IT 2107.12980, JHEP

Summary and outlook

- Leptogenesis: relation between neutrino physics and the very early Universe
- The baryon asymmetry can be produced for masses of right-handed neutrino ranging from ~ 0.1 GeV to GUT scale
- If the masses in the range $0.1 - 100$ GeV: experiment could reveal the origin of neutrino masses and the baryon asymmetry
- There are complementary search strategies for Heavy Neutral Leptons (LHC, SHiP, and FCC-ee)
- Heavy Neutral Leptons may hide even in what we think as “excluded” regions of the parameter space (140 MeV window, single mixing limits from LHC)

References

- Freeze-out of baryon number in low-scale leptogenesis
Shintaro Eijima, Mikhail Shaposhnikov, IT
[1709.07834](#), *JCAP* 11 (2017) 030
- Parameter space of baryogenesis in the ν MSM
Shintaro Eijima, Mikhail Shaposhnikov, IT
[1808.10833](#), *JHEP* 07 (2019) 077
- Uniting Low-Scale Leptogenesis Mechanisms
Juraj Klarić, Mikhail Shaposhnikov, IT
[2008.13771](#), *Phys.Rev.Lett.* 127 (2021) 11, 111802
- Reconciling resonant leptogenesis and baryogenesis via neutrino oscillations
Juraj Klarić, Mikhail Shaposhnikov, IT
[2103.16545](#), *Phys.Rev.D* 104 (2021) 5, 055010
- Dirac vs. Majorana HNLs (and their oscillations) at SHiP
Jean-Loup Tastet, IT
[1912.05520](#), *JHEP* 04 (2020) 005
- An allowed window for heavy neutral leptons below the kaon mass
Bondarenko et al.
[2101.09255](#) *JHEP* 07 (2021) 193
- Reinterpreting the ATLAS bounds on heavy neutral leptons in a realistic neutrino oscillation model
Jean-Loup Tastet, Oleg Ruchayskiy, IT
[2107.12980](#) *JHEP* 12 (2021) 182
- Heavy neutral leptons — Advancing into the PeV domain
Kevin Urquía-Calderón, IT, Oleg Ruchayskiy
[2206.04540](#) *JHEP* 08 (2023) 167

Backup slides

Description of low-scale leptogenesis

- Quantum kinetic equations (to capture HNL oscillations)

$$\begin{aligned}
 i \frac{dn_{\Delta_\alpha}}{dt} &= -2i \frac{\mu_\alpha}{T} \int \frac{d^3k}{(2\pi)^3} \text{Tr}[\Gamma_\alpha] f_N (1 - f_N) + i \int \frac{d^3k}{(2\pi)^3} \text{Tr}[\tilde{\Gamma}_\alpha (\delta\bar{\rho}_N - \delta\rho_N)], \\
 i \frac{d\delta\rho_N}{dt} &= -i \frac{d\rho_N^{eq}}{dt} + [H_N, \rho_N] - \frac{i}{2} \{\Gamma, \delta\rho_N\} - \frac{i}{2} \sum_\alpha \tilde{\Gamma}_\alpha \left[2 \frac{\mu_\alpha}{T} f_N (1 - f_N) \right], \\
 i \frac{d\delta\bar{\rho}_N}{dt} &= -i \frac{d\rho_N^{eq}}{dt} - [H_N, \bar{\rho}_N] - \frac{i}{2} \{\Gamma, \delta\bar{\rho}_N\} + \frac{i}{2} \sum_\alpha \tilde{\Gamma}_\alpha \left[2 \frac{\mu_\alpha}{T} f_N (1 - f_N) \right].
 \end{aligned}$$

Not affected by
sphalerons

$$n_{\Delta_\alpha} = L_\alpha - B/3$$

Susceptibility matrix –
spectator effects

$$\mu_\beta = \omega_{\beta\alpha} n_{\Delta_\alpha}$$

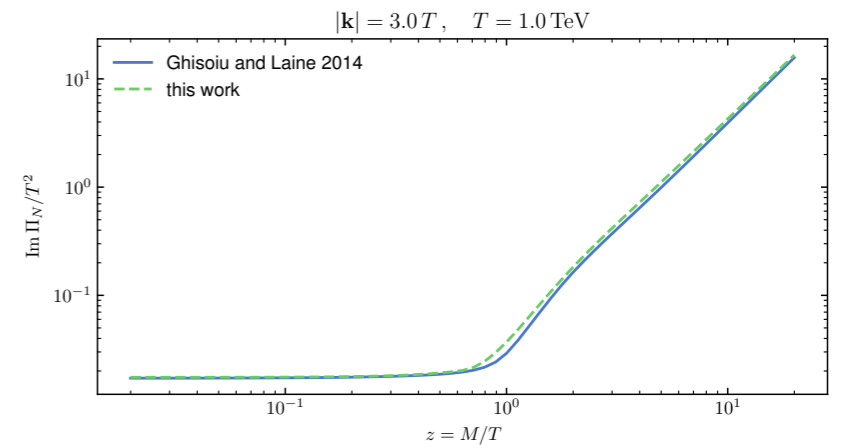
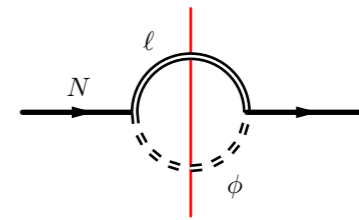
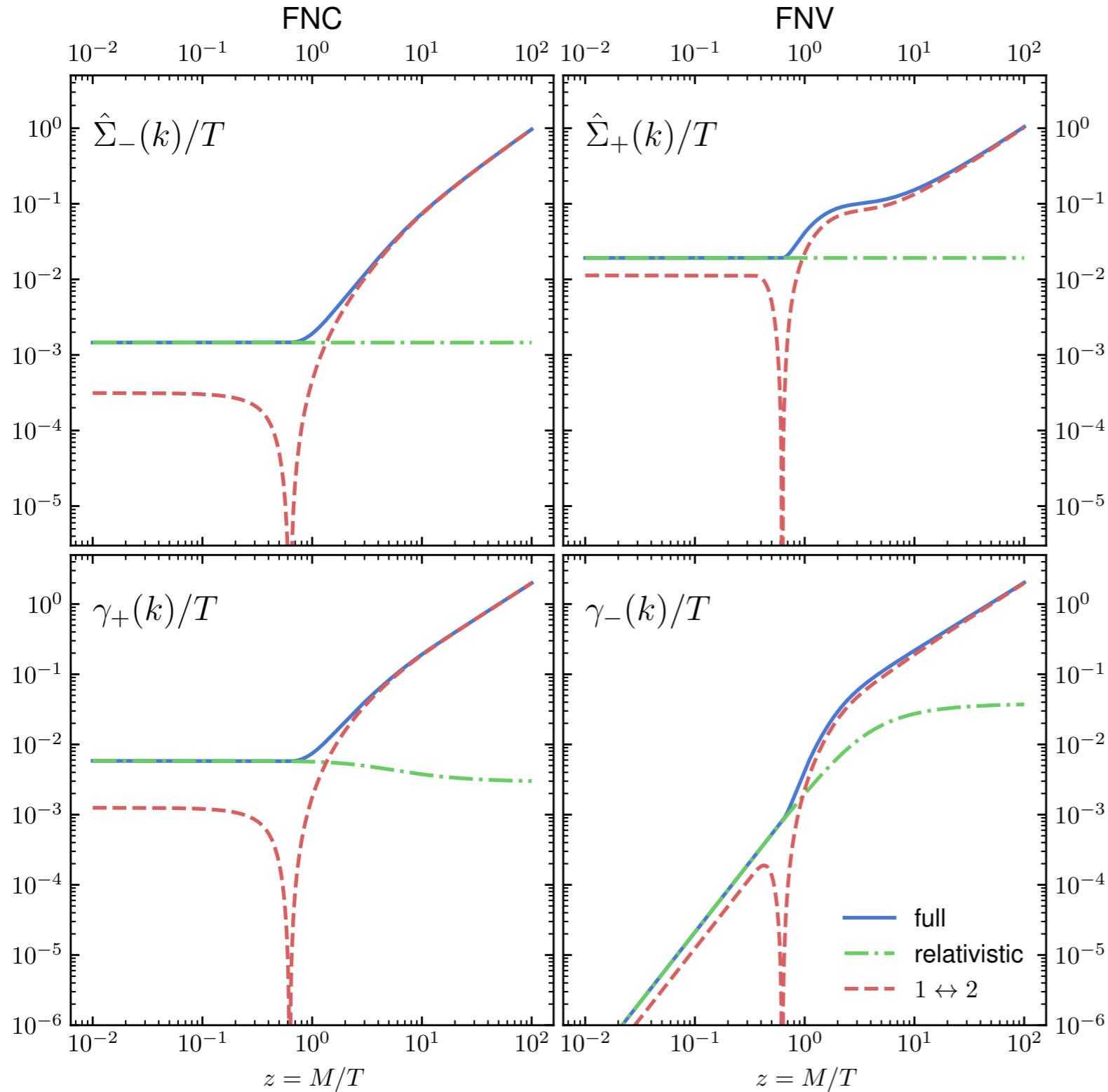
2x2 HNL matrix of
densities

$$\rho_N \quad \delta\rho_N = \rho_N - \rho_N^{eq}$$

- The equations must be solved numerically
- Scan over 6-dimensional parameter space (mass of N, mass splitting, phases of Yukawas)

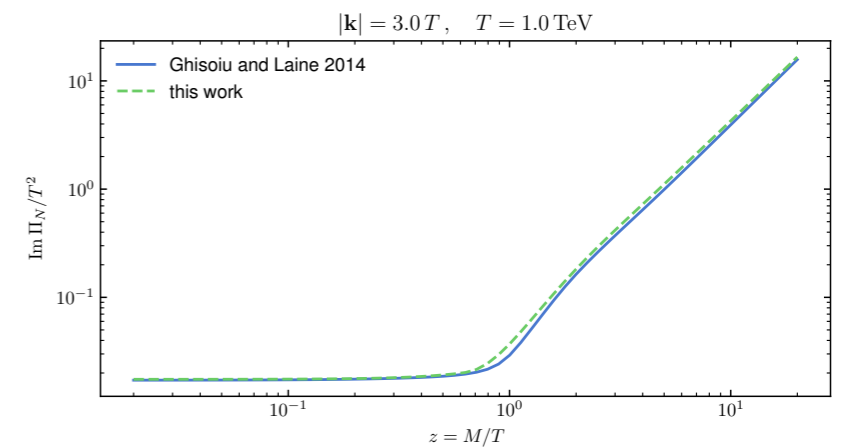
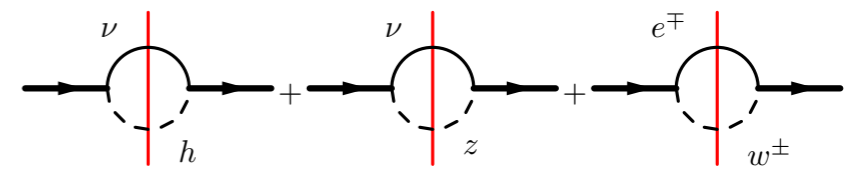
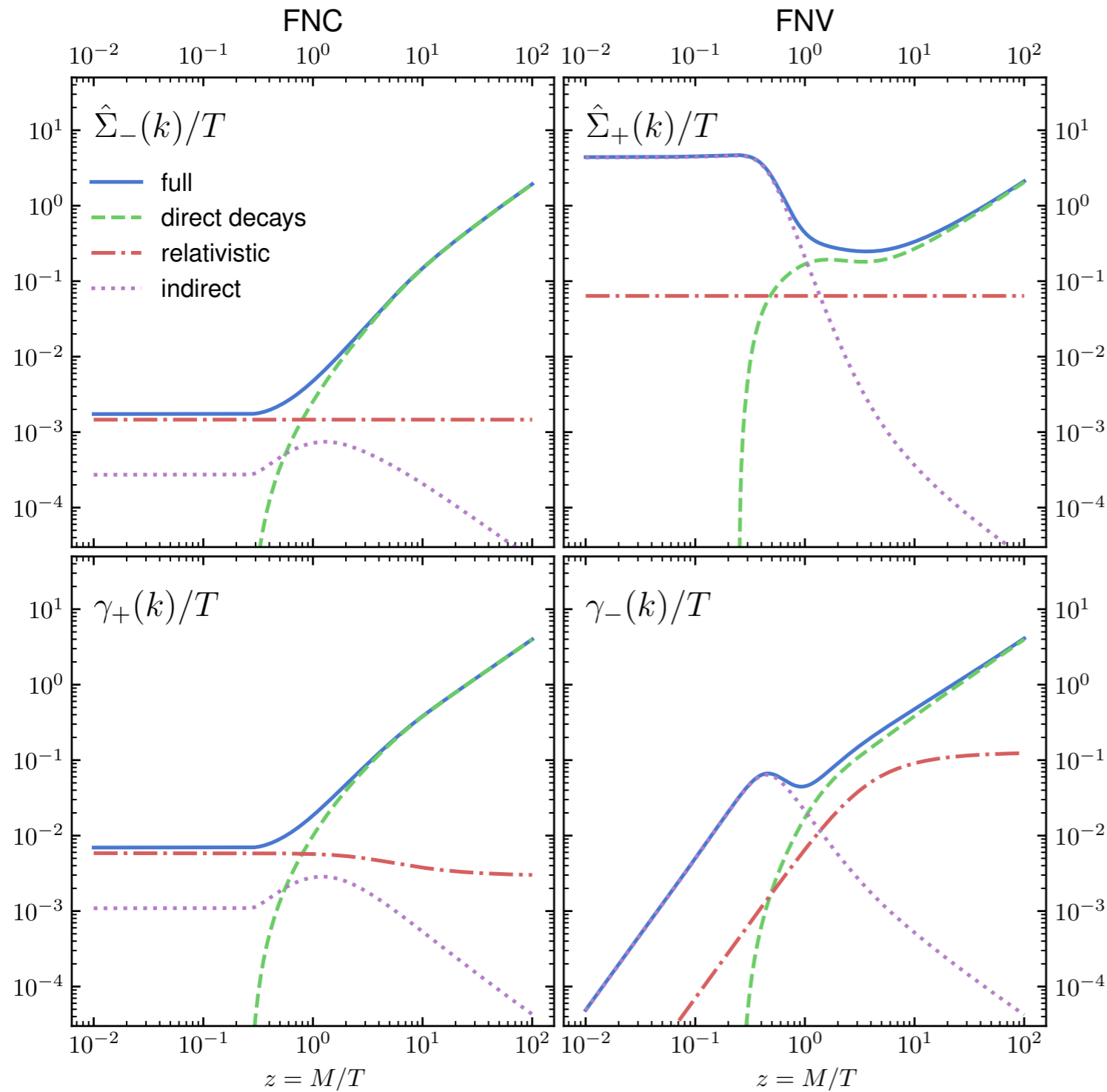
The rates

$$|\mathbf{k}| = 3.0T, \quad T = 1.0 \text{ TeV}$$

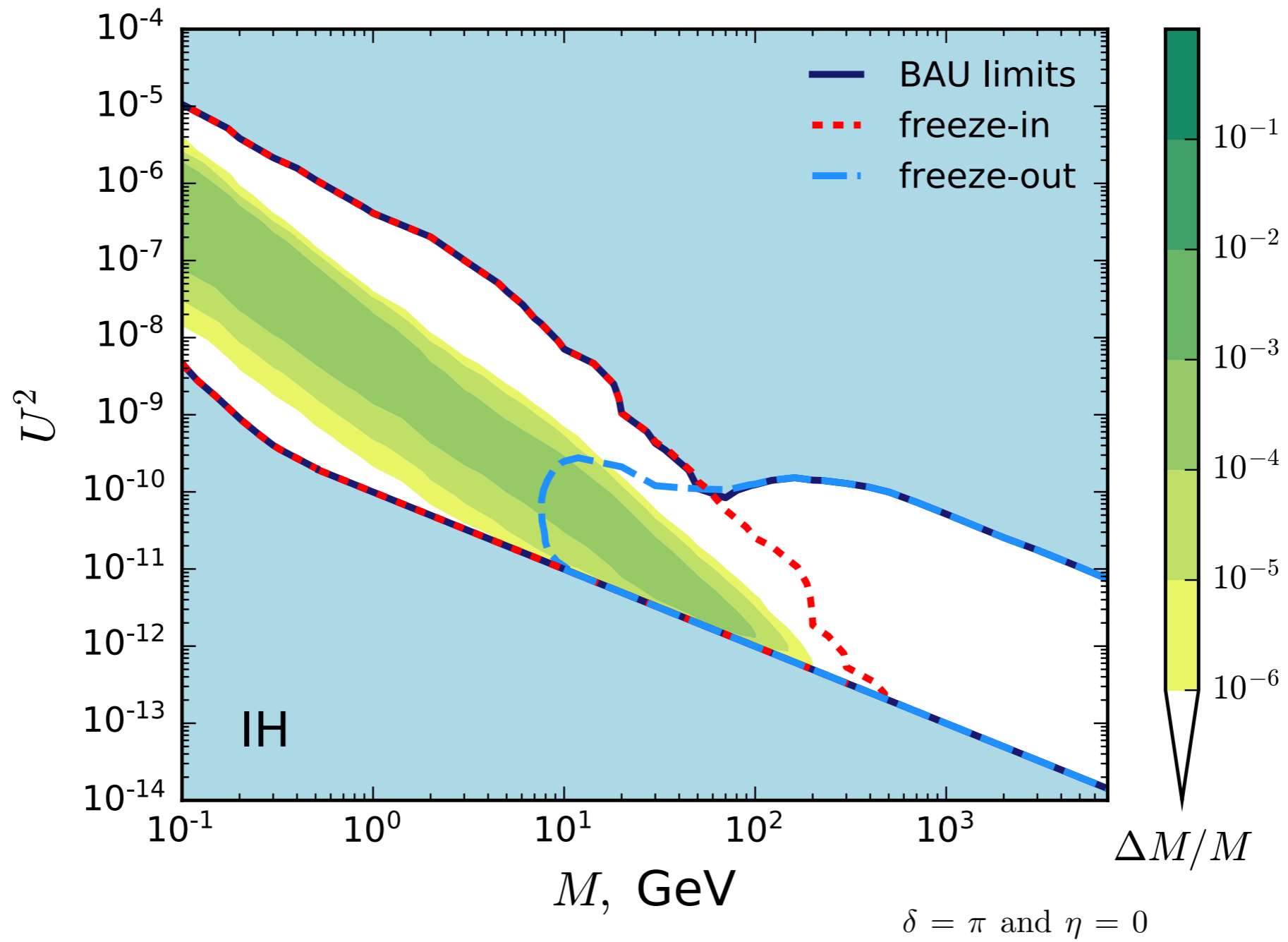


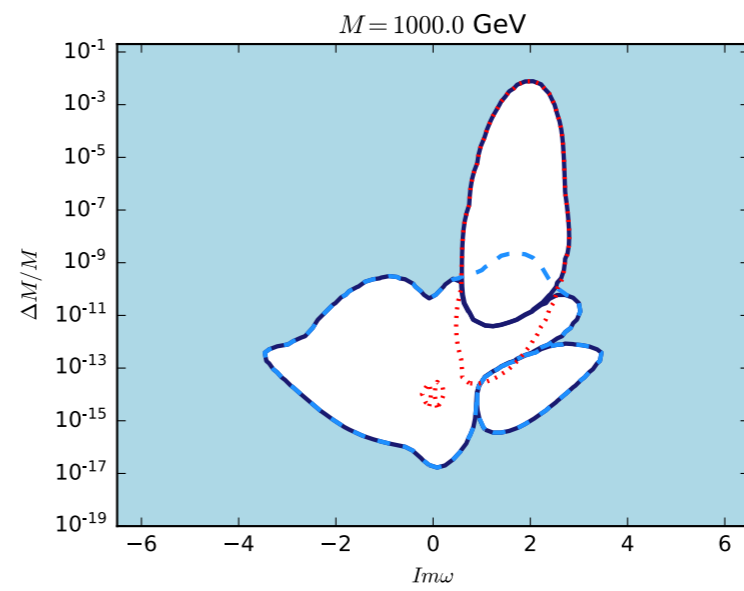
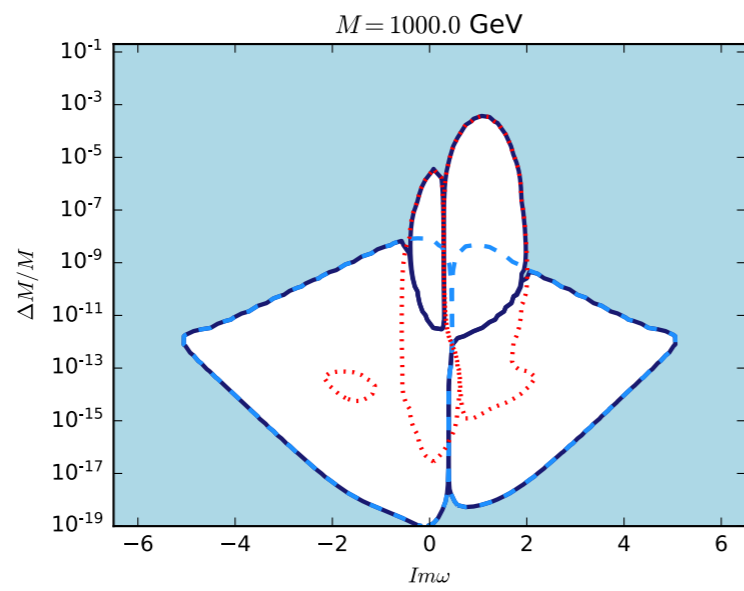
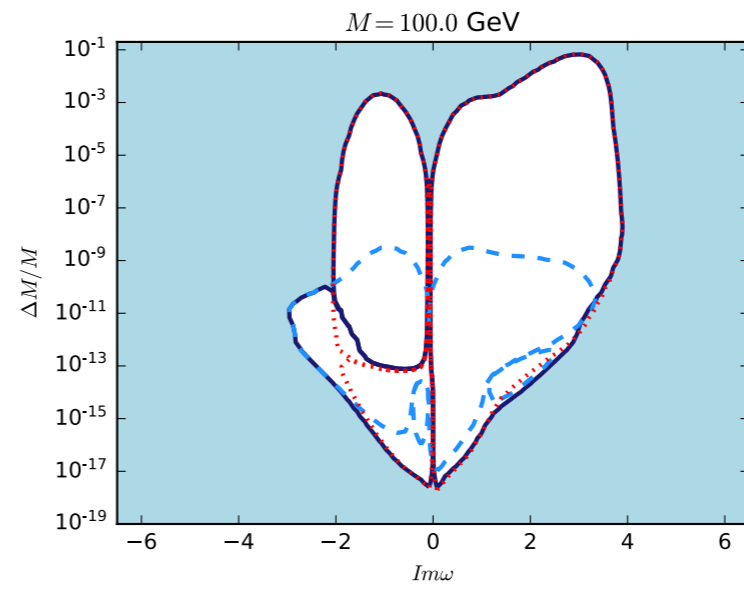
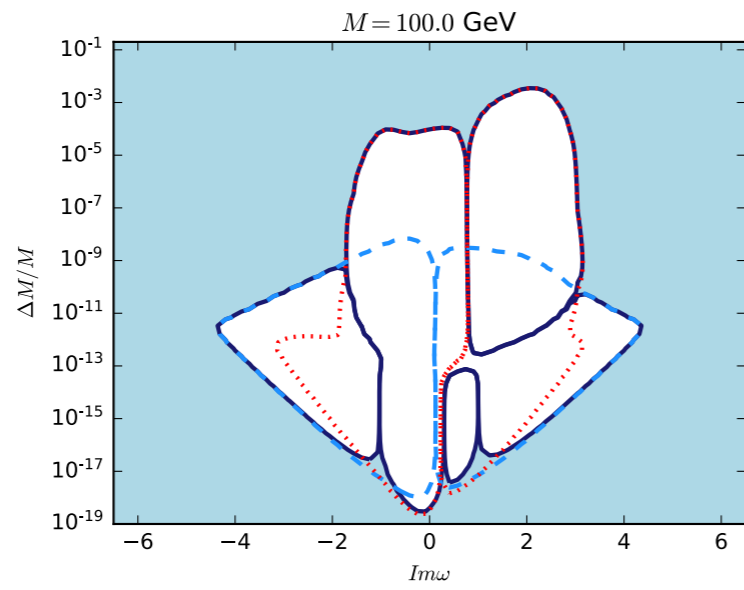
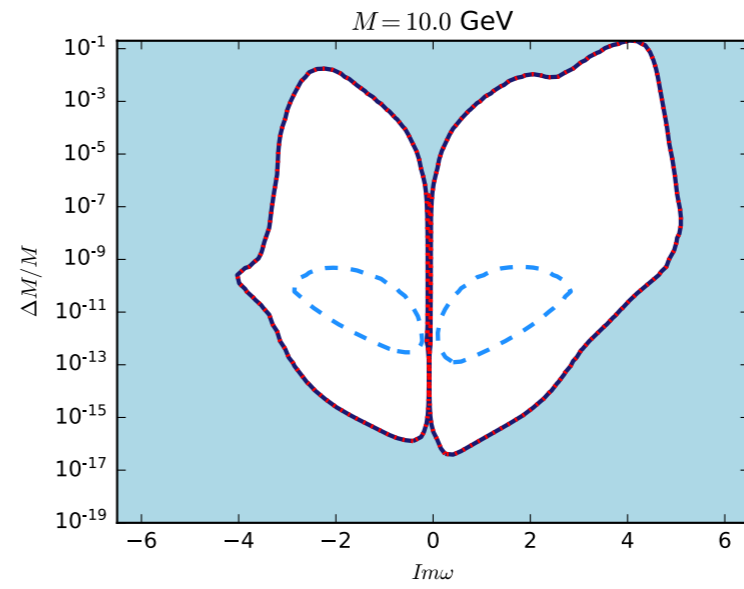
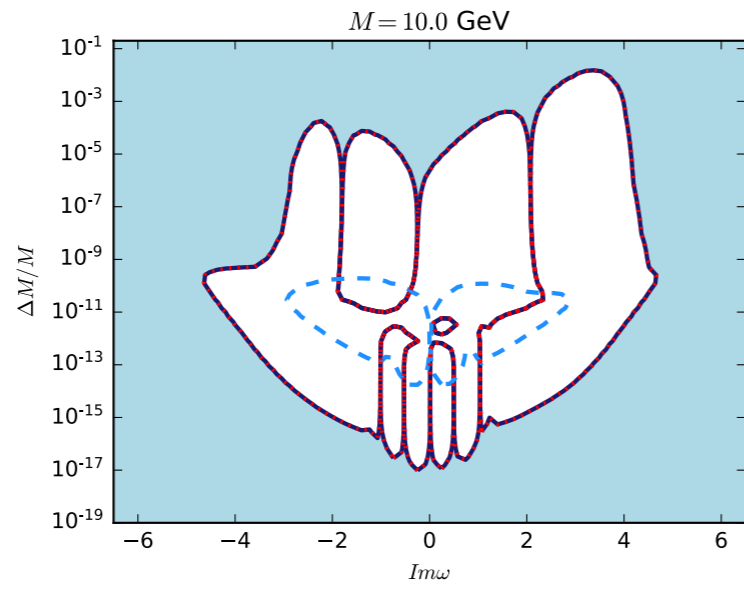
The rates

$$|\mathbf{k}| = 3.0T, \quad T = 140 \text{ GeV}$$



flavour hierarchical washout





Casas-Ibarra parametrization

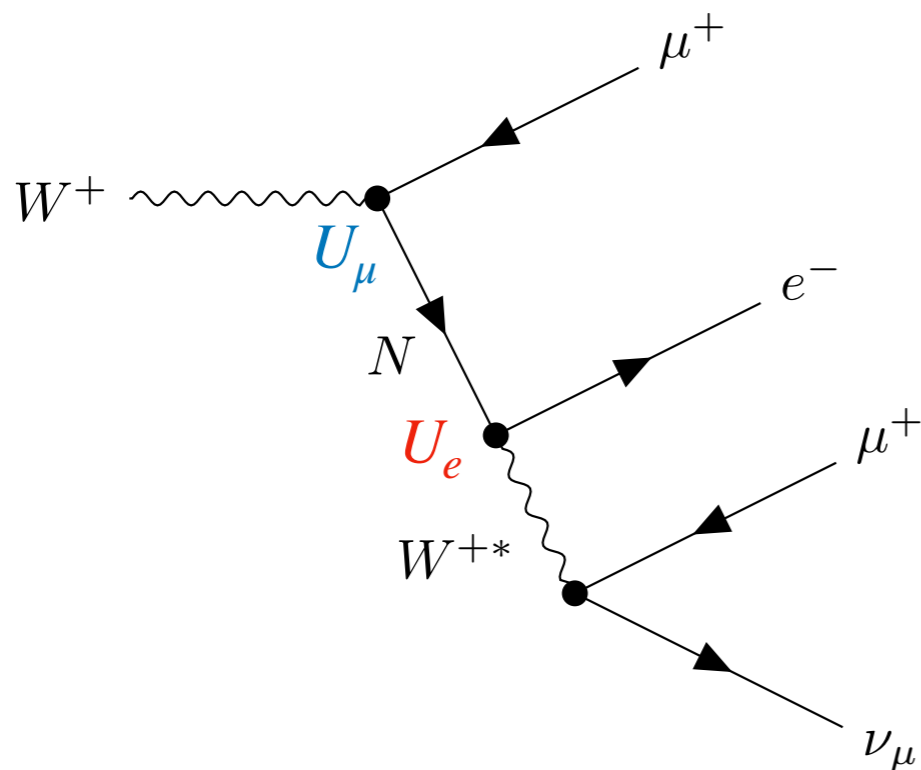
$$F = \frac{i}{v} U_\nu \sqrt{m_\nu^{\text{diag}}} \mathcal{R} \sqrt{M_M} ;$$

$$\mathcal{R}^{\text{NH}} = \begin{pmatrix} 0 & 0 \\ \cos \omega & \sin \omega \\ -\xi \sin \omega & \xi \cos \omega \end{pmatrix}, \quad \mathcal{R}^{\text{IH}} = \begin{pmatrix} \cos \omega & \sin \omega \\ -\xi \sin \omega & \xi \cos \omega \\ 0 & 0 \end{pmatrix}$$

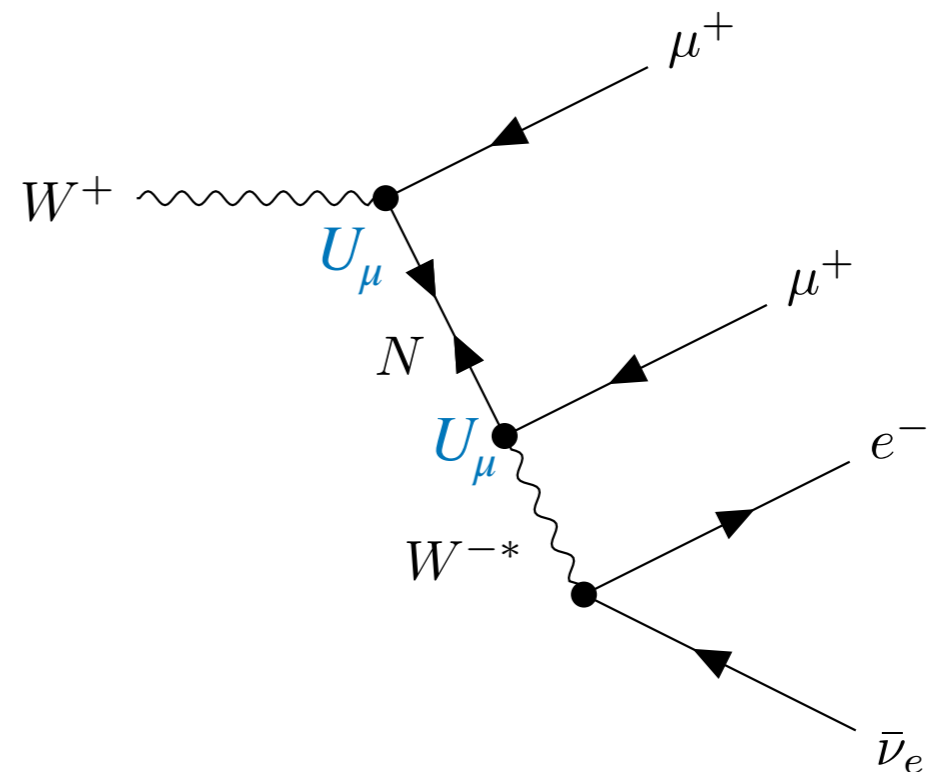
$M, \text{ GeV}$	$\log_{10}(\Delta M/M)$	$\text{Im } \omega$	$\text{Re } \omega$	δ	η
[0.1 – 7000]	[-19, -0.5]	[-7, 7]	[0, π]	[0, 2π]	[0, 2π]

Neutrino oscillation data and reinterpretation

ATLAS triplepton search [1905.09787](https://arxiv.org/abs/1905.09787)



(a) LNC



(b) LNV

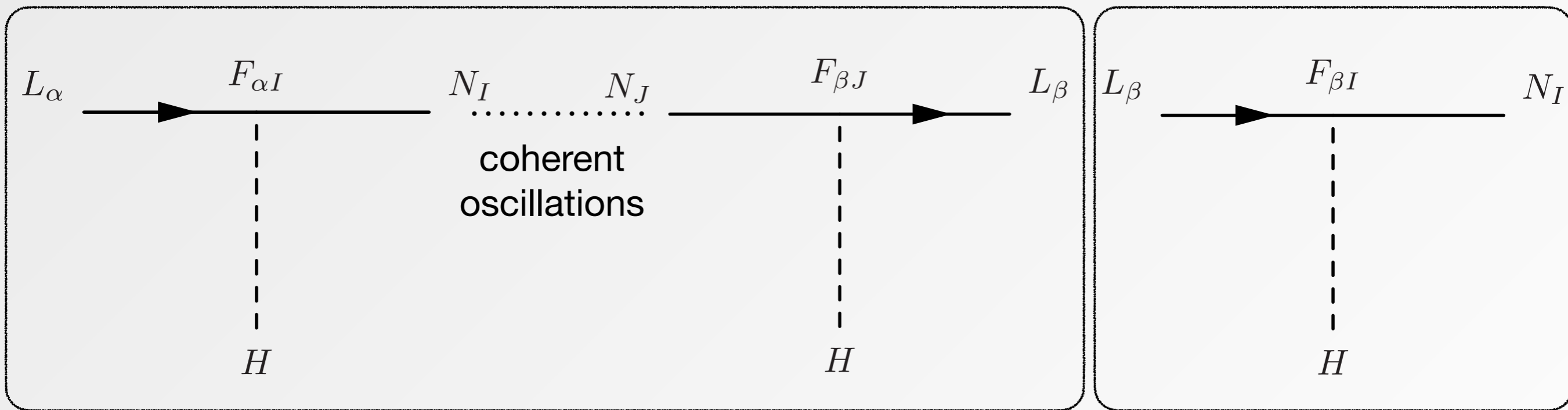
LNC cannot be probed under single mixing assumption

Thanks to Jean-Loup and Oleg
ATLAS now considers different mixing patterns!
<https://arxiv.org/abs/2204.11988>

$$\mathcal{L} = \mathcal{L}_{SM} + i\bar{\nu}_{R_I}\gamma^\mu\partial_\mu\nu_{R_I} - F_{\alpha I}\bar{L}_\alpha\tilde{\Phi}\nu_{R_I} - \frac{M_{IJ}}{2}\bar{\nu}_{R_I}^c\nu_{R_J} + h.c.$$

BAU generation

time \longrightarrow



No lepton asymmetry

SM species

are in equilibrium

$L \rightarrow N$ is out of equilibrium

Individual lepton asymmetries.

$$n_{L_\alpha} \neq n_{\bar{L}_\alpha}$$

Total lepton asymmetry

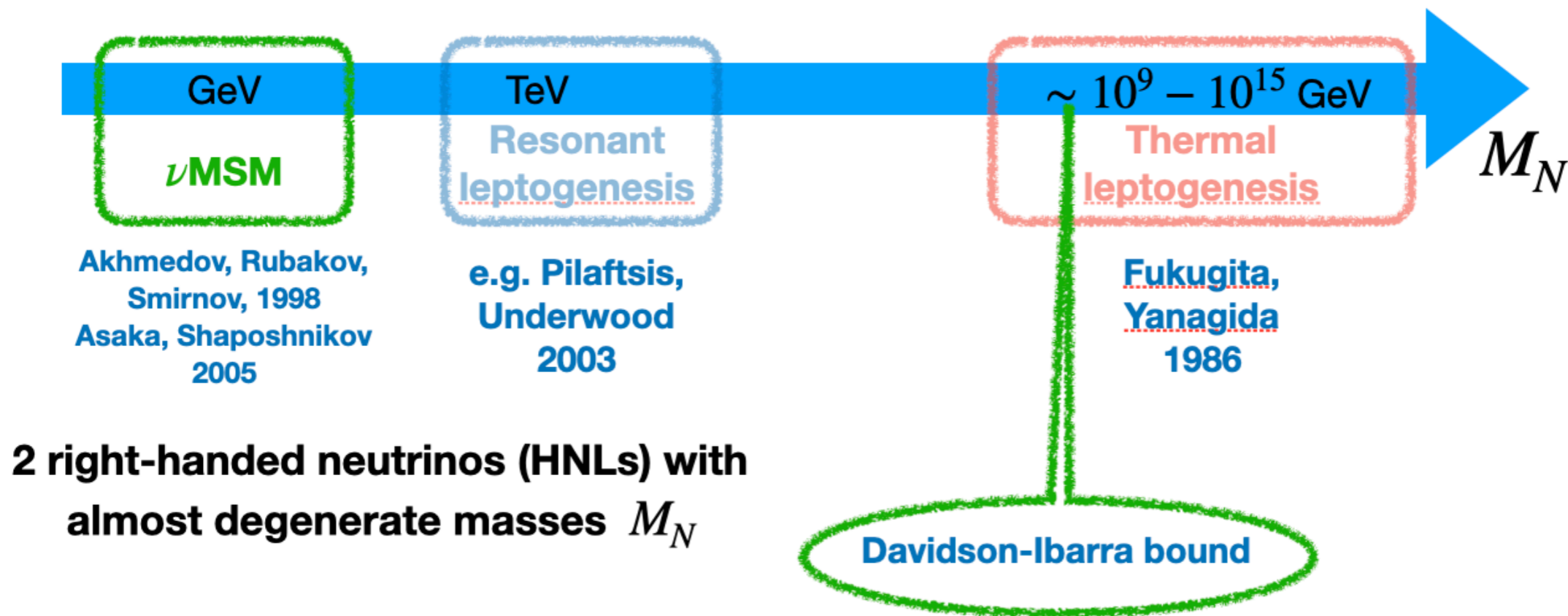
$$\Gamma(L_\alpha \rightarrow L_\beta) \neq \Gamma(\bar{L}_\alpha \rightarrow \bar{L}_\beta)$$

Leptogenesis

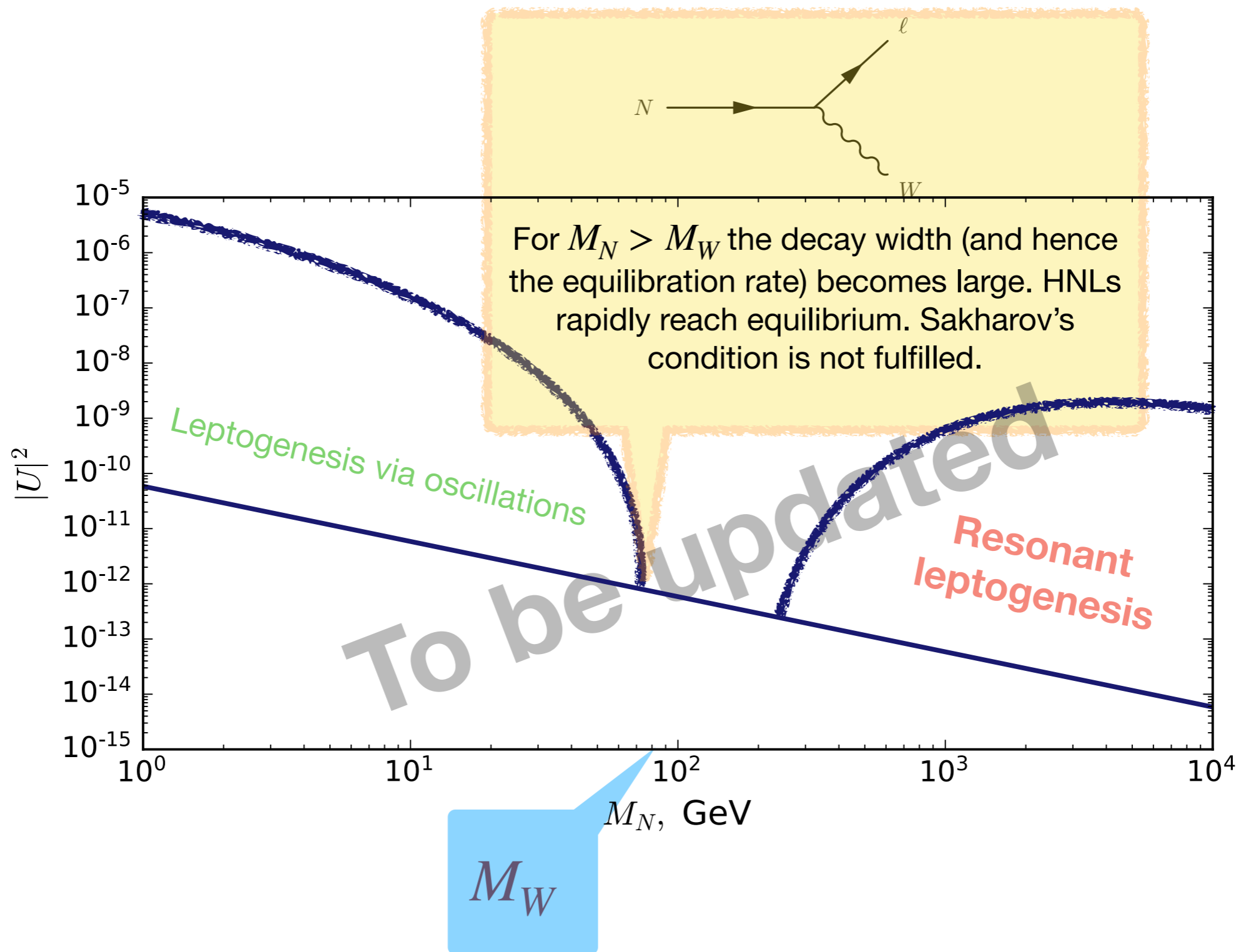
$$\eta \equiv \frac{n_B}{n_\gamma} \simeq 6.2 \times 10^{-10}$$

baryon asymmetry from lepton asymmetry
by the sphaleron processes

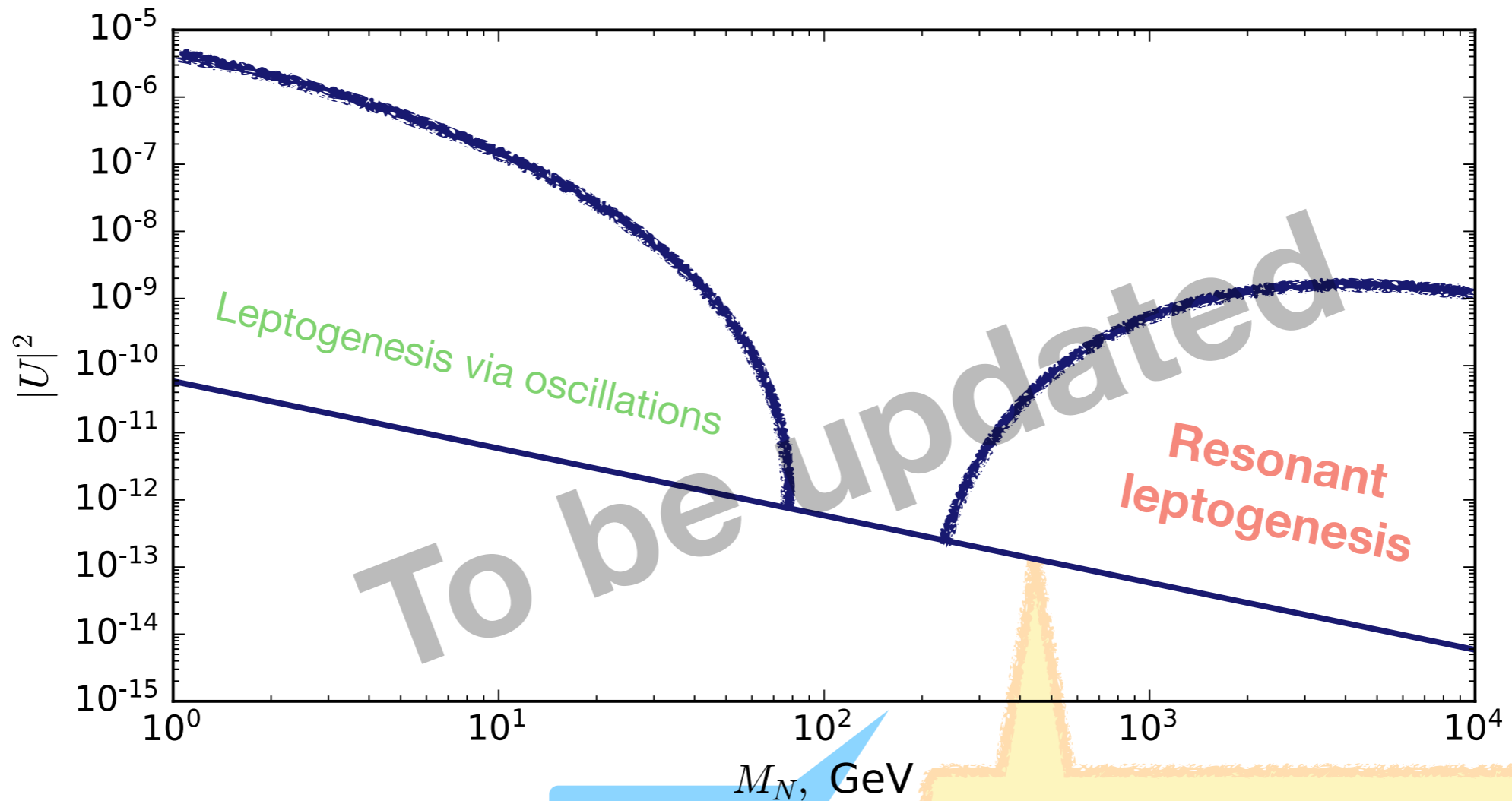
Kuzmin, Rubakov
Shaposhnikov 1985



Different leptogenesis mechanisms?



Different leptogenesis mechanisms?

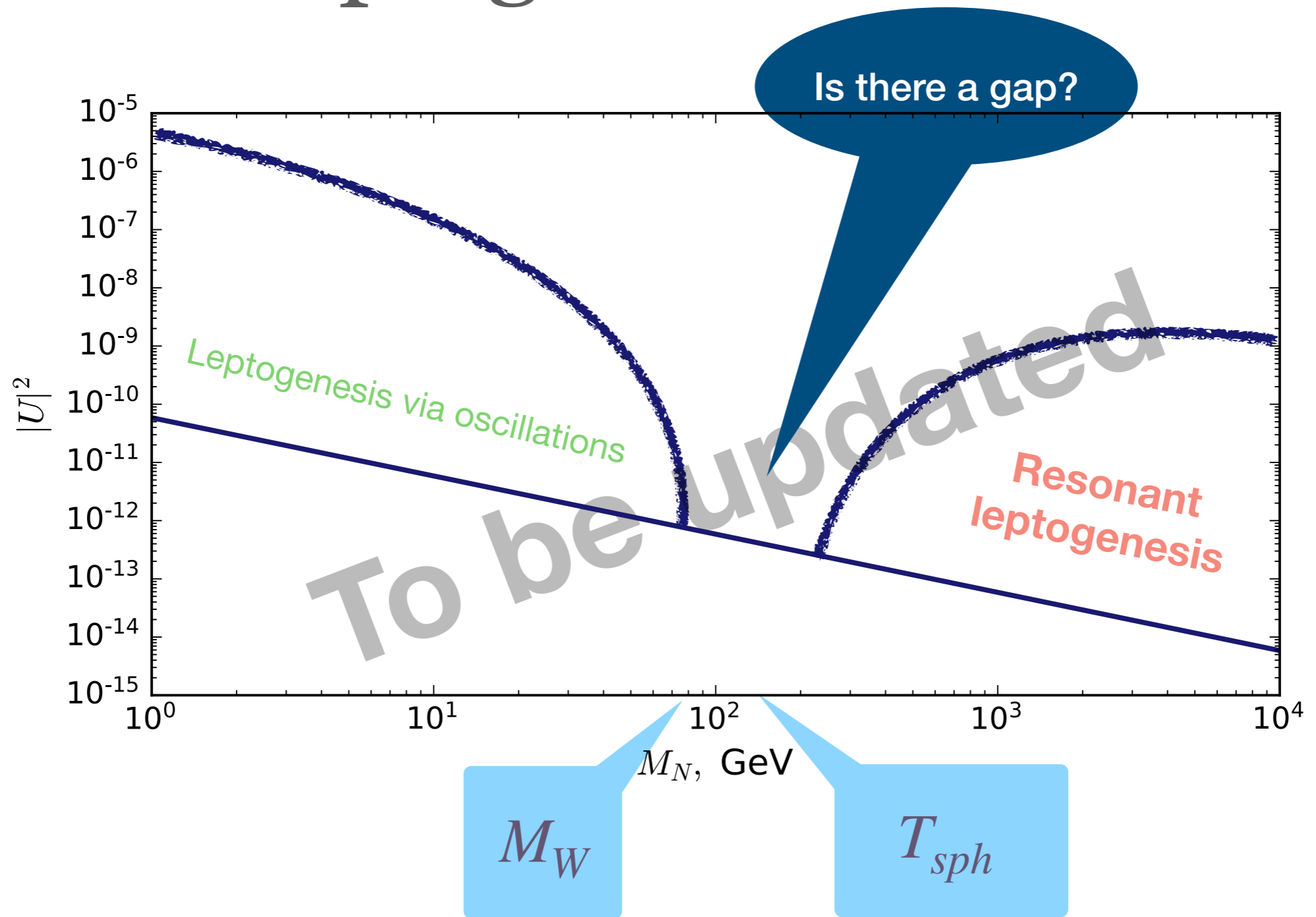


T_{sph}

Lepton asymmetry is generated in decays of N

If $M_N < 130 \text{ GeV}$ one can expect that this asymmetry is not transferred to BAU

Different leptogenesis mechanisms?



$$d\Gamma_{\alpha\beta}^{\text{LNC/LNV}}(\tau) \cong 2 |\Theta_{\alpha 1}|^2 |\Theta_{\beta 1}|^2 \left(1 \pm \cos(\Delta M \tau)\right) e^{-\Gamma \tau} d\hat{\Gamma}_{\alpha\beta}^{\text{LNC/LNV}}$$

$\Delta M \tau \ll 2\pi$ (Dirac-like limit)

$\Delta M \tau \gg 2\pi$ (Majorana-like limit)

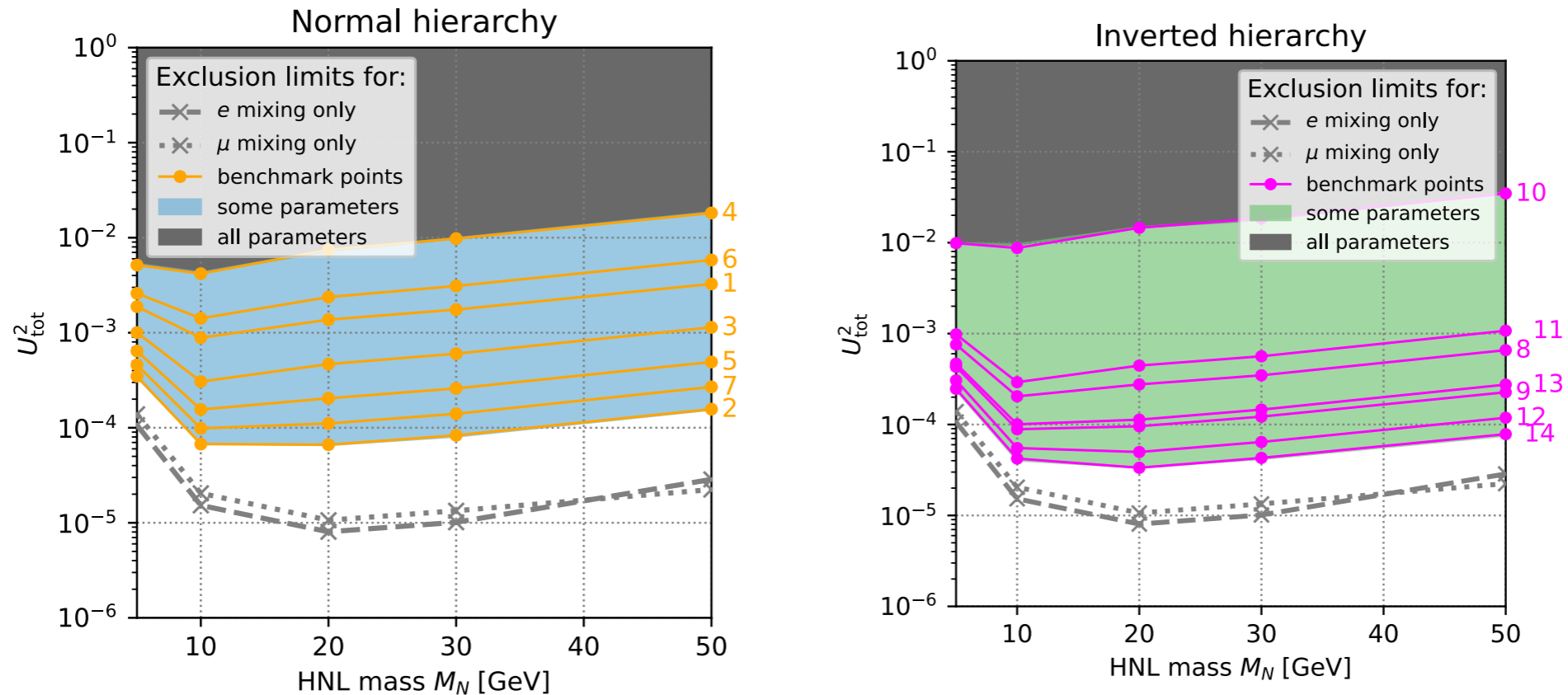
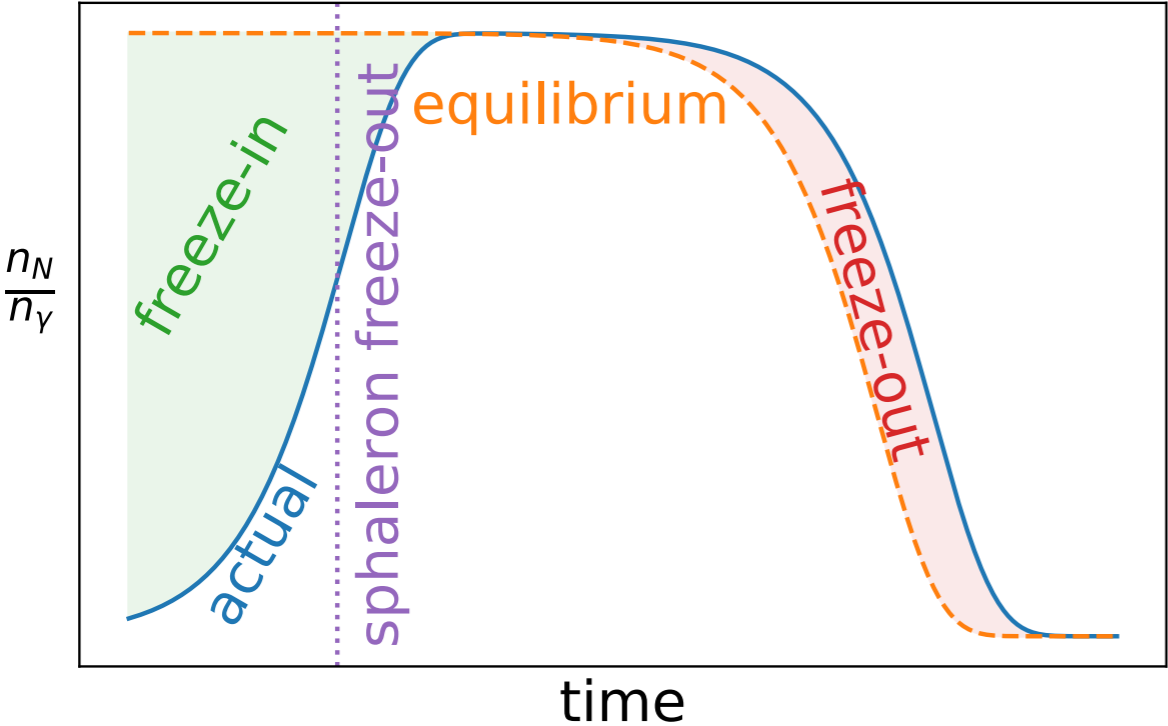
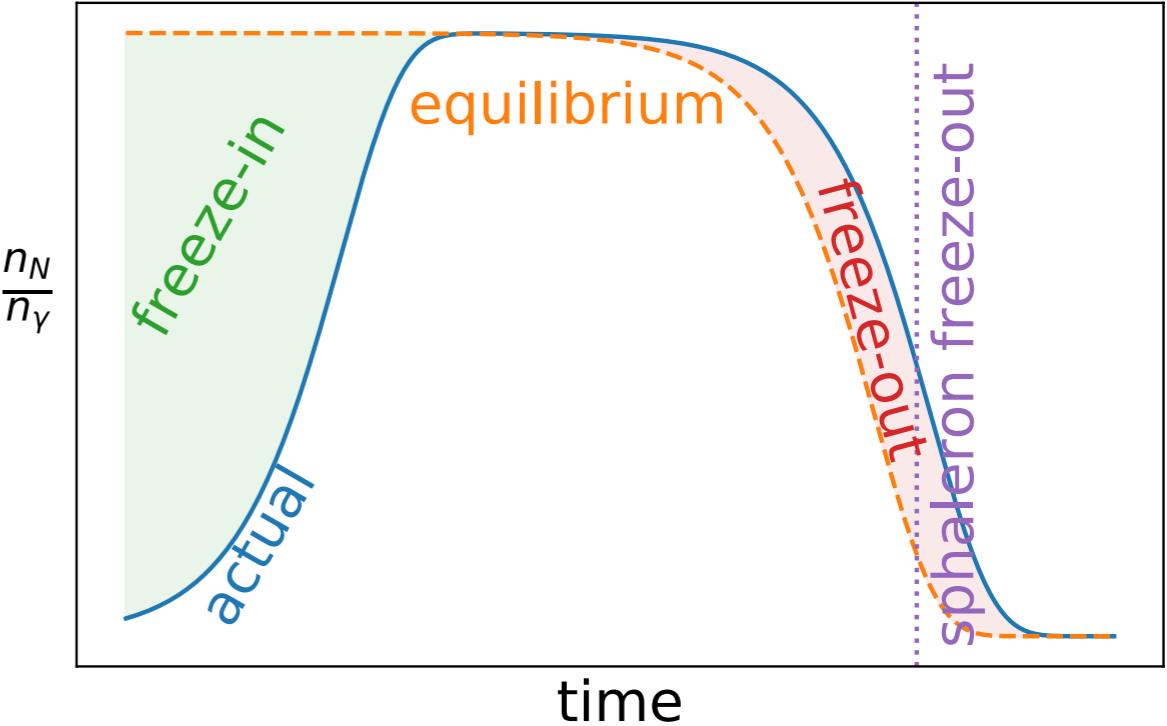


Figure 8: Same as figure 7, but for a **Dirac-like** HNL pair. The single-flavor mixing limits are grayed out because this search has *no sensitivity* to the Dirac-like case under this assumption; instead the limits for the Majorana-like case are given for comparison.

More accurate classification of Leptogenesis mechanisms



“Leptogenesis via oscillations”



“Resonant Leptogenesis”